

# Scalable quantum computing with neutral atom qubits

Mark Saffman



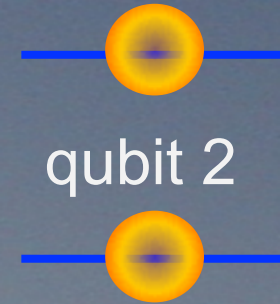
- Neutral atom qubits

- Efficient quantum search

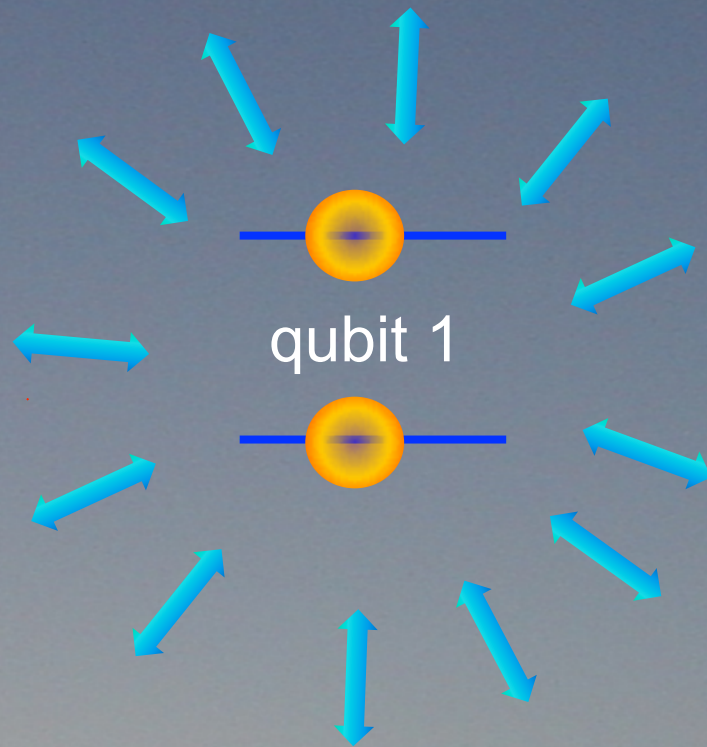
- Experimental status



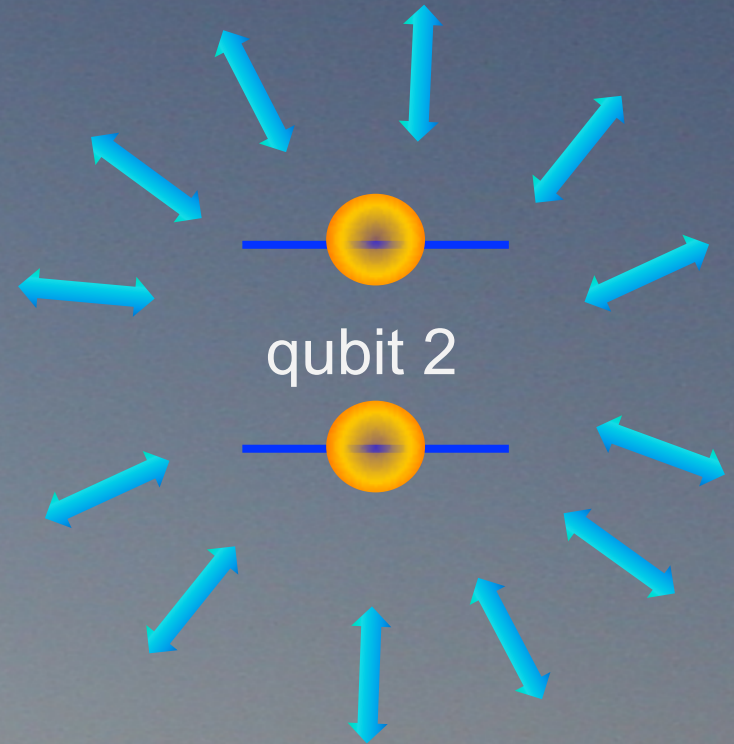
# The challenge of quantum computing



# The challenge of quantum computing

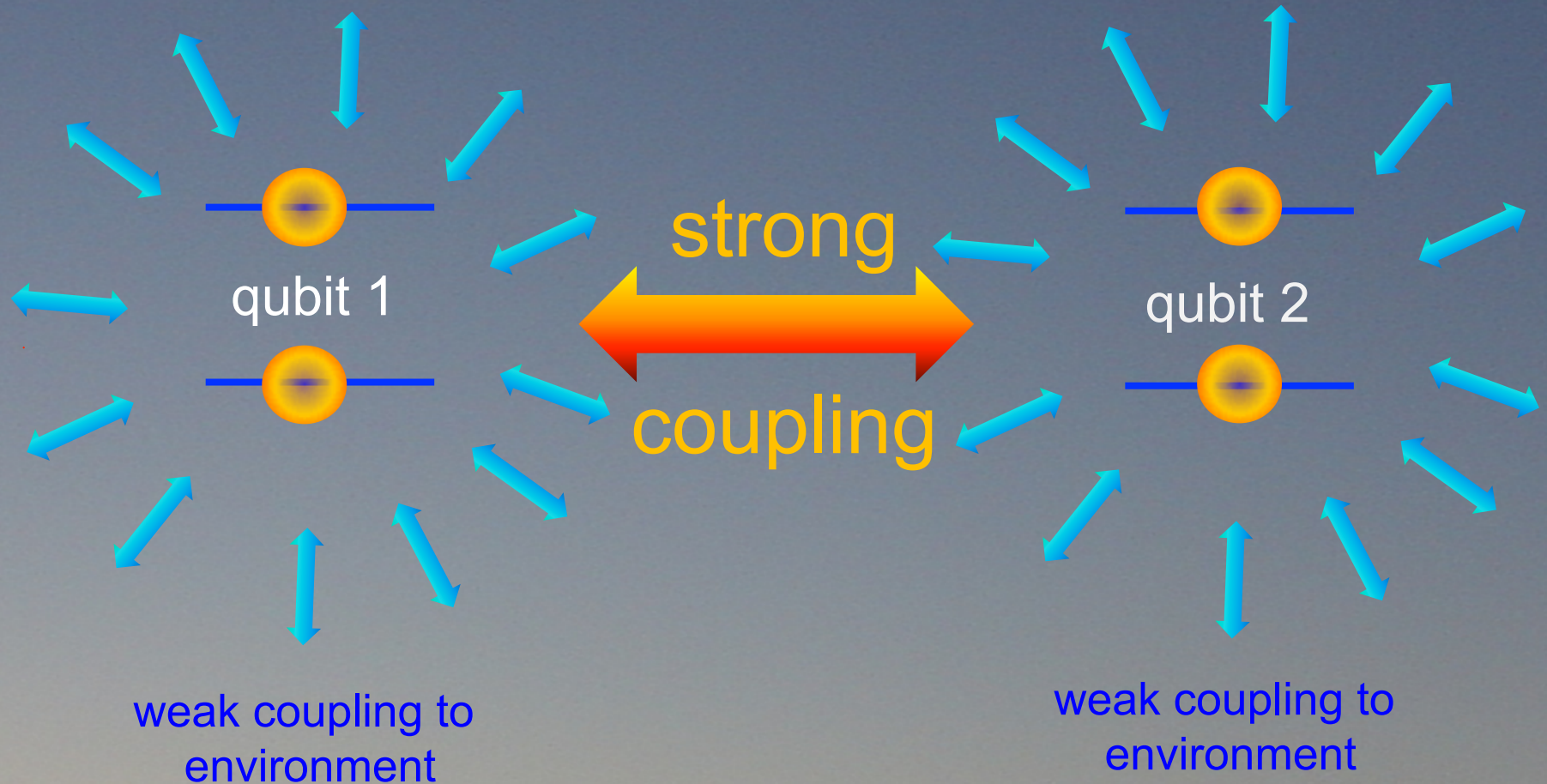


weak coupling to  
environment

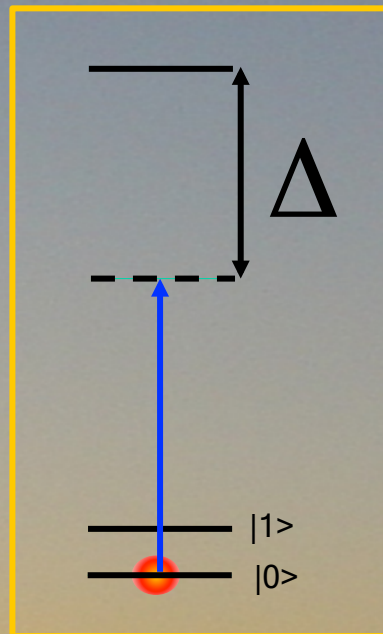
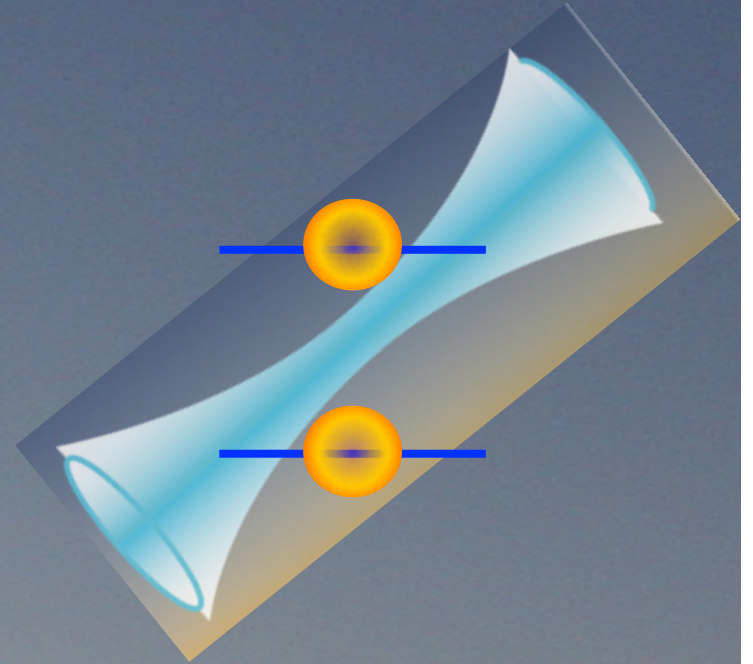
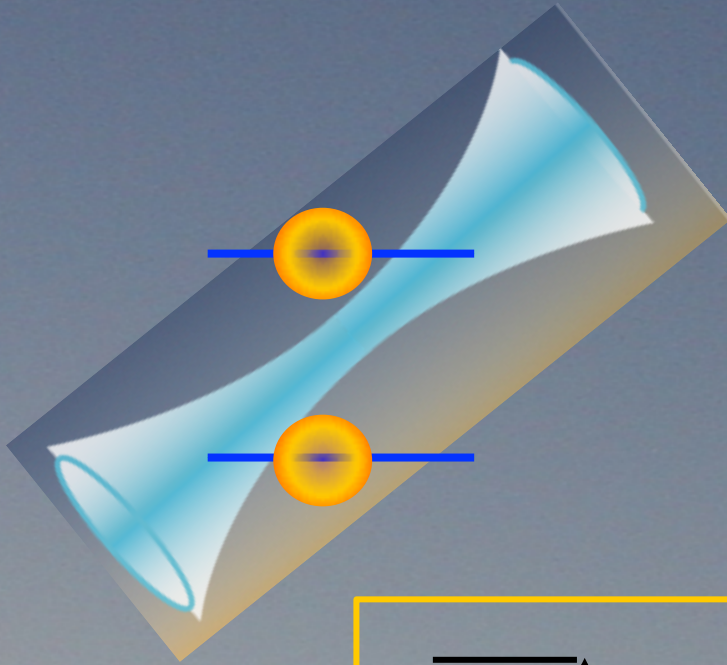


weak coupling to  
environment

# The challenge of quantum computing



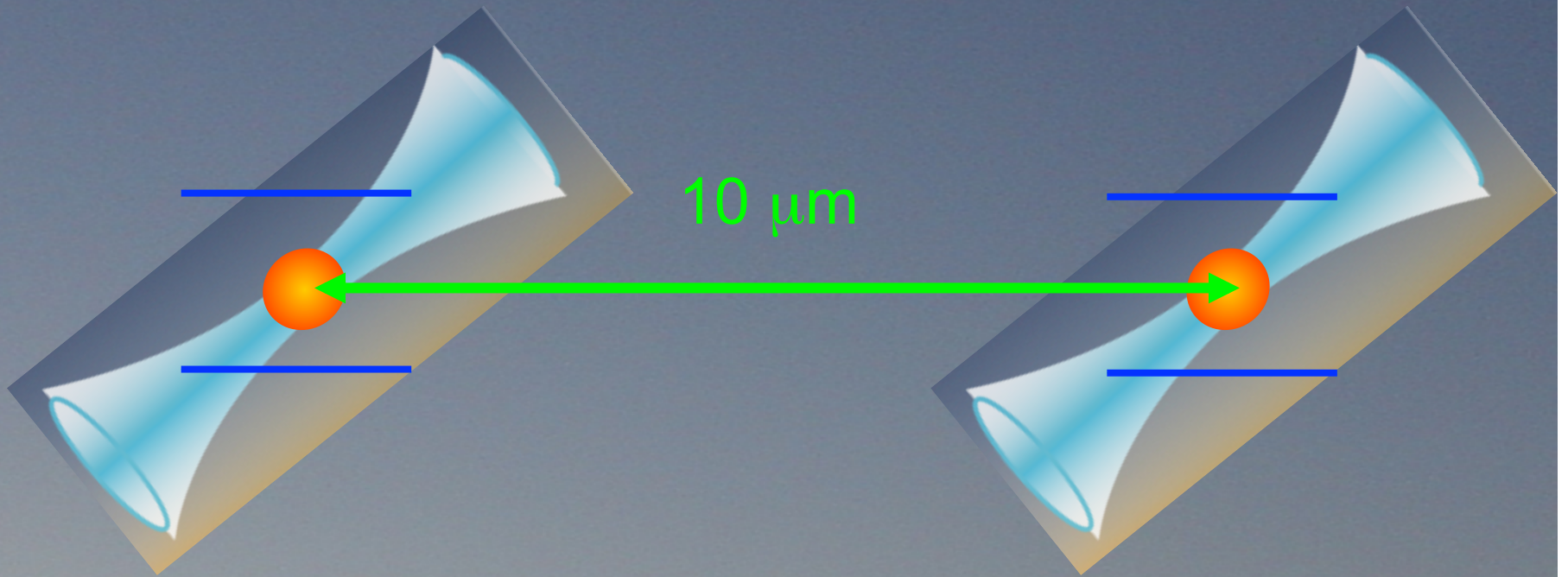
# Neutral atoms in optical traps



Trap depth  $\sim 1/\Delta$

Decoherence rate  $\sim 1/\Delta^2$

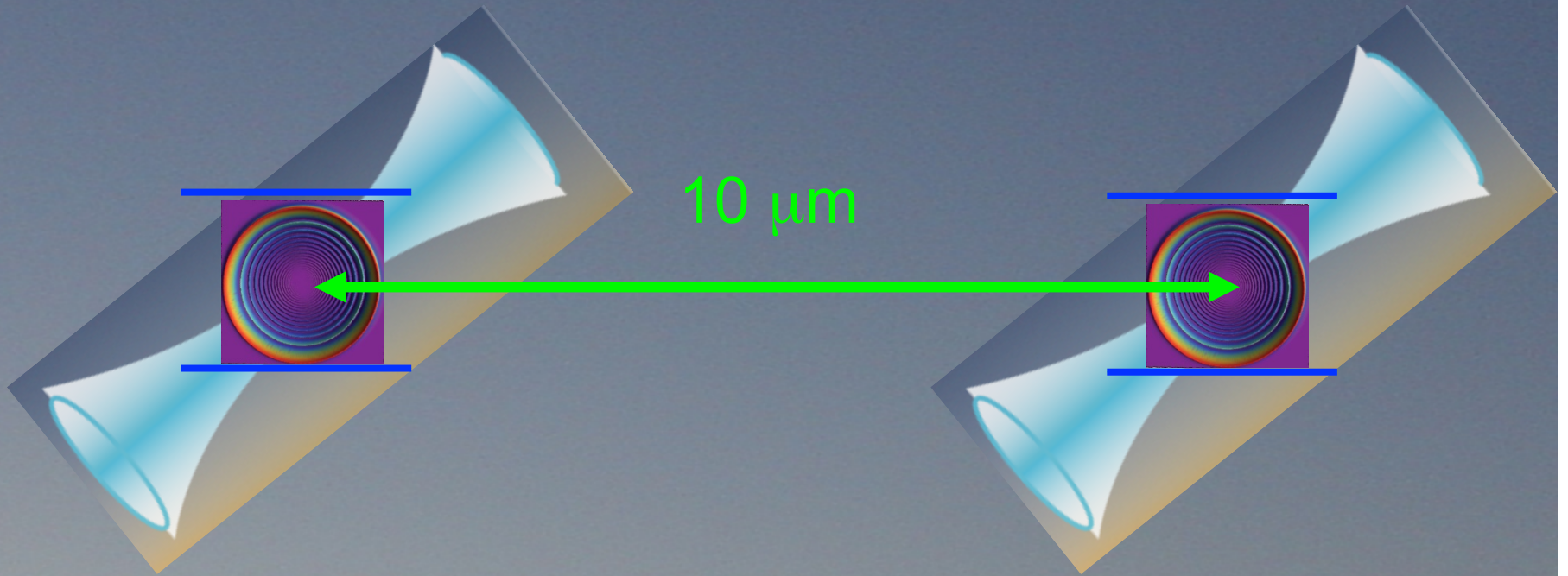
# Long range interactions



Rb-Rb ground state  
magnetostatic interaction

$$\Delta E \sim 100 \mu\text{Hz}$$

# Long range interactions



Rb-Rb ground state  
magnetostatic interaction

$$\Delta E \sim 100 \mu\text{Hz}$$

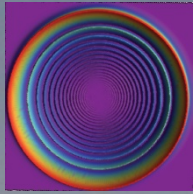
Rydberg  $n=100$   
van der Waals interaction

$$\Delta E \sim 100 \text{MHz}$$

**12 orders of magnitude!**

# Rydberg atoms

Highly excited atoms with exaggerated properties.

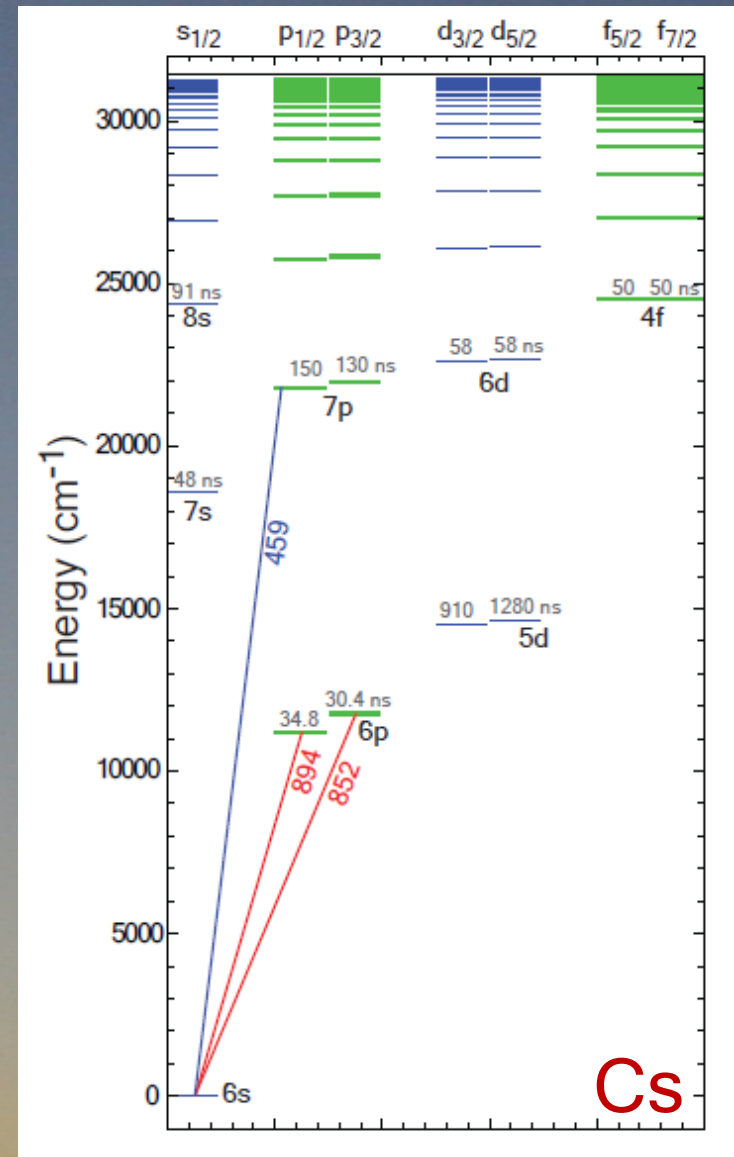


$$\langle r \rangle \sim a_0 n^2$$

$$V_{vdw} \sim n^{12}$$

$$\tau \sim n^3$$

$$\text{Fidelity of } C_Z \sim 1 - \frac{a}{(V\tau)^{2/3}}$$





# Fast Quantum Gates for Neutral Atoms

Rydberg  
Blockade

D. Jaksch, J.I. Cirac, and P. Zoller

*Institut für Theoretische Physik, Universität Innsbruck, Technikerstrasse 25, A-6020 Innsbruck, Austria*

S.L. Rolston

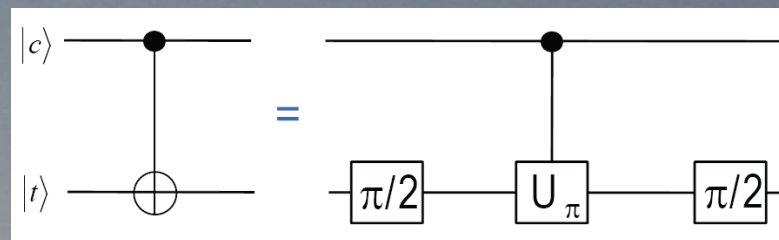
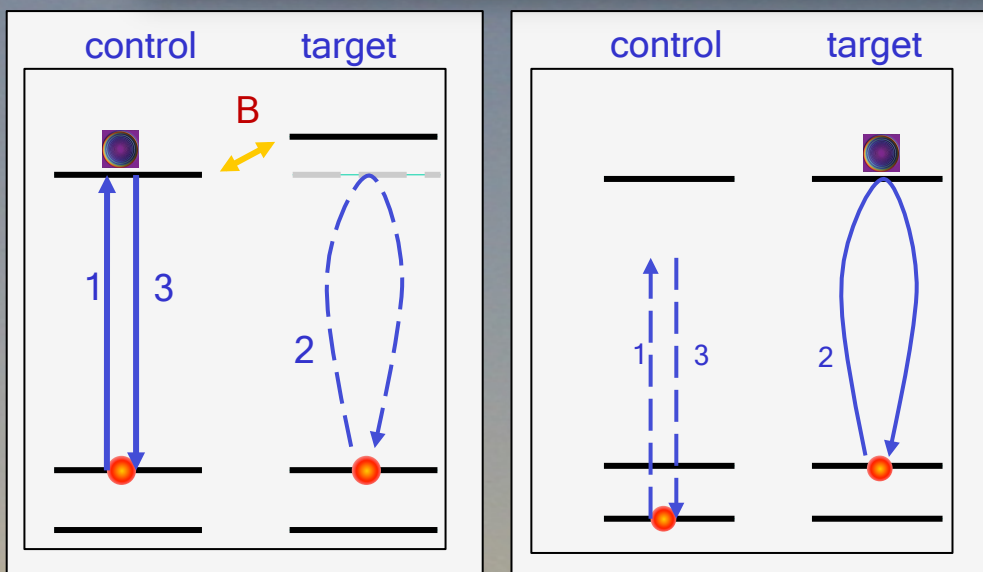
*National Institute of Standards and Technology, Gaithersburg, Maryland 20899*

R. Côté<sup>1</sup> and M.D. Lukin<sup>2</sup>

<sup>1</sup>*Physics Department, University of Connecticut, 2152 Hillside Road, Storrs, Connecticut 06269-3046*

<sup>2</sup>*ITAMP, Harvard-Smithsonian Center for Astrophysics, Cambridge, Massachusetts 02138*

(Received 7 April 2000)



conditional phase of  $\pi$  on target

# Rydberg atom quantum information related experiments

## Quantum information with Rydberg atoms

M. Saffman and T. G. Walker

*Department of Physics, University of Wisconsin, 1150 University Avenue, Madison, Wisconsin 53706, USA*

K. Molmer

*Lundbeck Foundation Theoretical Center for Quantum System Research, Department of Physics and Astronomy, University of Aarhus, DK-8000 Århus C, Denmark*



- UW Madison
- University of Connecticut
- University of Michigan
- University of Virginia
- University of Oklahoma
- Penn State
- MIT
- Georgia Tech
- Sandia National Lab



- Waterloo



- University of Heidelberg
- University of Stuttgart
- Max Planck Institute for Quantum Optics, Garching
- Technical University of Darmstadt



- ENS Paris
- Institut d'Optique, Palaiseau
- Laboratoire Aimé Cotton, Université Paris Sud



- ETH, Zurich



- University of Durham
- Open University, Milton Keynes



- University of Amsterdam



- University of Pisa



- RAS Novosibirsk



- University of Electro-Communications



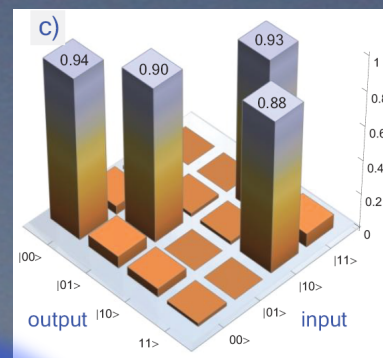
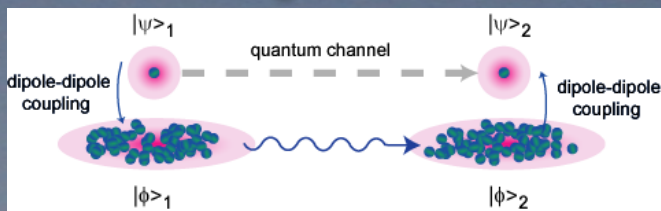
- Shanxi University



- São Paulo

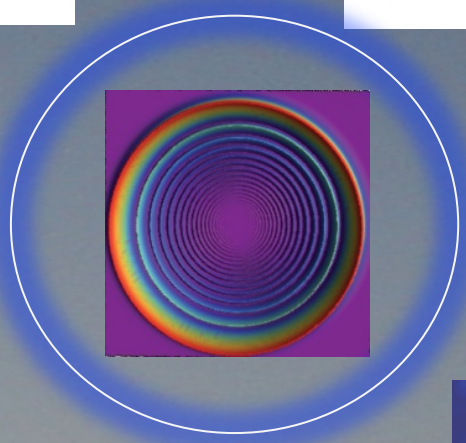
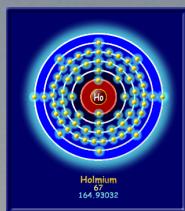
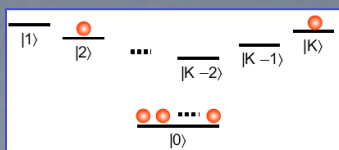
# Quantum Information with Rydberg atoms

## Ensembles for quantum networking

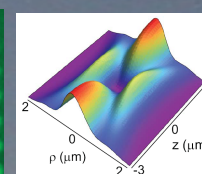
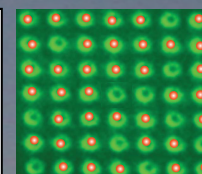
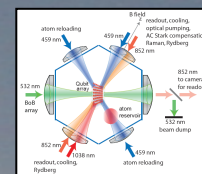


Quantum gates, entanglement

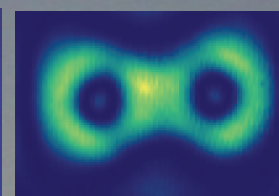
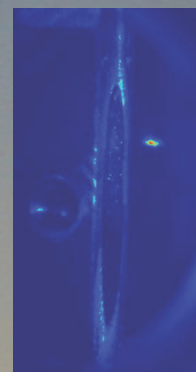
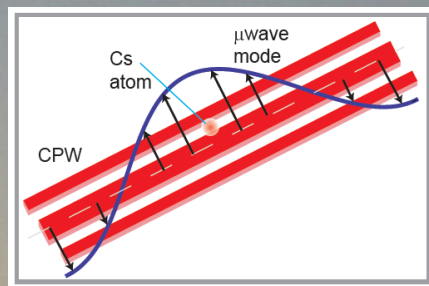
No collective encoding



## AQuA - atomic Qubit array



hybrid interfaces

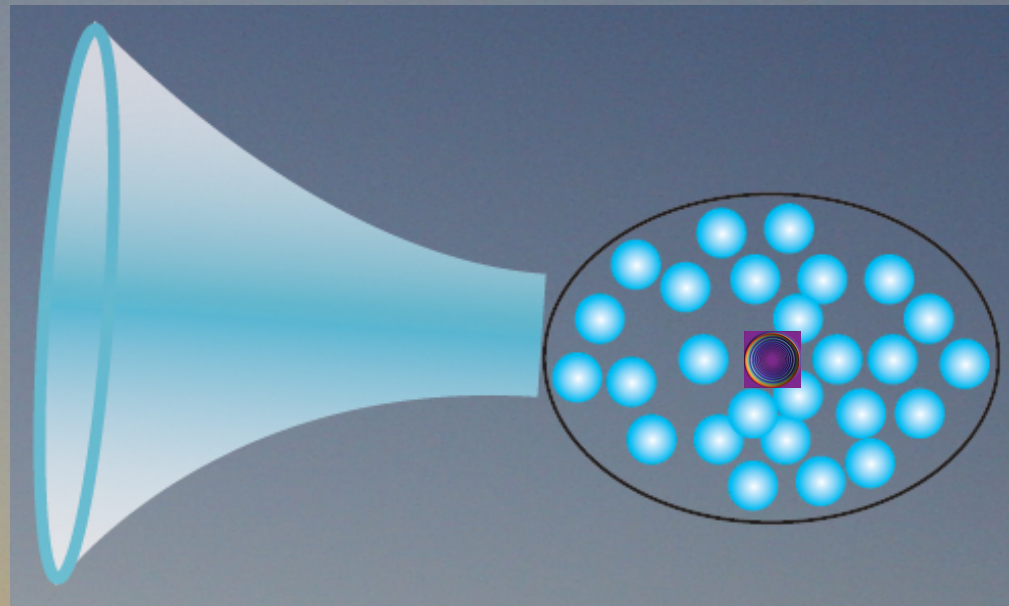


Holographic chip traps  
atom - surface interactions

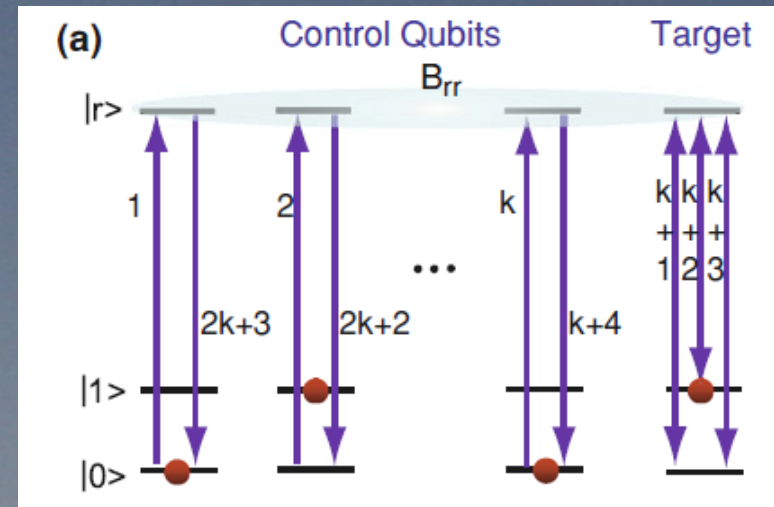
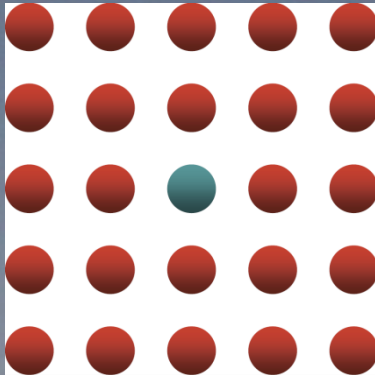
# Efficient quantum search

- As with classical computing we can use universal quantum hardware to run generic algorithms.
- Nevertheless some hardware is better suited for certain tasks, (e.g. graphics accelerators,..).
- Rydberg blockade enables significant speedup for the quantum search algorithm.
- Ensemble blockade

One atom can control the state of  $N$  other atoms inside a blockade sphere.



# Blockade and multi-bit gates



- $k$  control atoms and 1 target atom
- If one or more control atoms is in state 0, there will be a Rydberg excitation which blocks target atom pulses
- $C_k$ NOT gate needs only  $2k+3$  laser pulses
- Best known qu. circuit:  $32k-120$  gates,  $k \geq 5$ .

**Table 1** Sequential addressing  $C_k$ NOT gate errors from Eq. (2) averaged over a square lattice for several different Rydberg Cs  $ns$  levels

$n$	$\tau$ ( $\mu$ s)	$d$ ( $\mu$ m)	$k = 3$	$k = 8$	$k = 15$	$k = 24$	$k = 35$
50	63	1	0.003	0.04	0.30	>1	>1
75	170	1	0.0003	0.0017	0.0071	0.026	0.078
100	330	2.13	0.0009	0.0021	0.010	0.035	0.11
125	540	3.65	0.0003	0.0024	0.011	0.039	0.12
150	820	5.66	0.0004	0.0028	0.012	0.042	0.13

For  $n = 75$  the Rabi frequency was set to  $\Omega/2\pi = (45, 29, 20, 15, 12)$  MHz for  $k = (3, 8, 15, 24, 35)$

$2^{35} \sim 34$  billion

# Grover search

- Register

$$\sum_x c_x |x\rangle$$

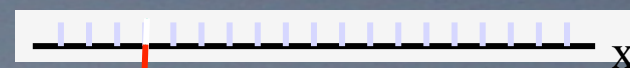


- Wish to find marked element

$$|x_0\rangle$$

- 1) Sign change

$$c_{x_0} \rightarrow -c_{x_0}$$



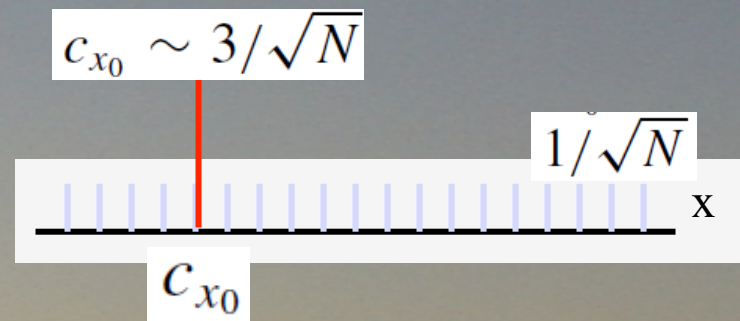
- 2) Inversion about mean

$$\bar{c} = \frac{1}{N} \sum_x c_x$$

$$c_x \rightarrow \frac{1}{N} \sum_{x'} c_{x'} - (c_x - c_{x'})$$

- Factor 3 enhancement after first step.

- Succeeds with high probability after  $N^{1/2}$  steps vs.  $N/2$  for classical search.

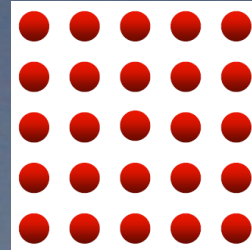


# Multi-bit gates and Grover search

## Efficient Grover search with Rydberg blockade

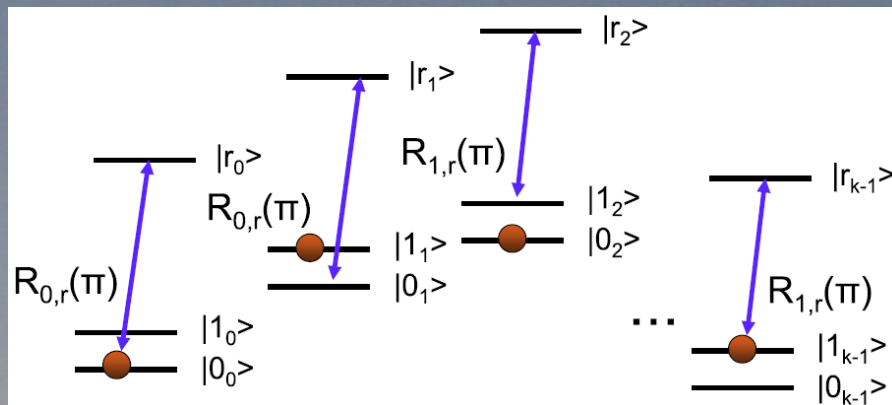
Klaus Mølmer<sup>1,3</sup>, Larry Isenhower<sup>2</sup> and Mark Saffman<sup>2</sup>

$k = \log_2(N)$  qubits



sign change

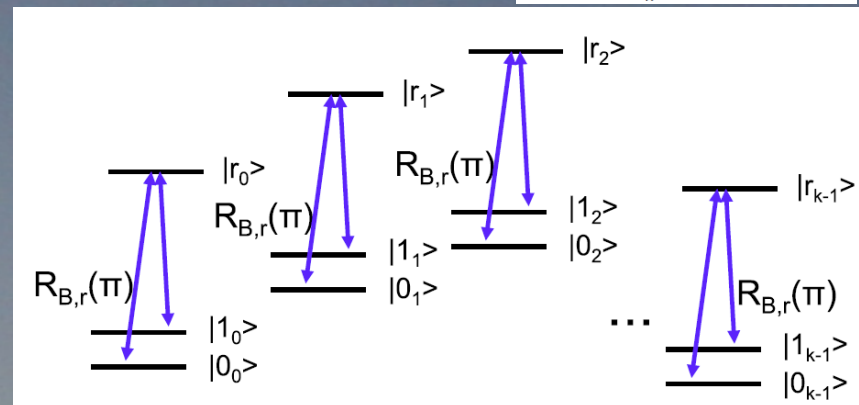
$$c_{x_0} \rightarrow -c_{x_0}$$



**Figure 2.** Grover conditional phase: successive  $\pi$ -pulse excitation transfer of all atoms from one of their qubit states to the Rydberg state. The lasers couple the states  $|1 - b_i\rangle$ , complementary to the bit value of the marked element  $x_0$ , to the Rydberg states. If one atom is already excited, no further excitation occurs, and hence quantum register components with more atoms populating the states coupled to the laser field will only be excited once, while the state with no atoms coupled to the laser fields will not be excited at all. A second set of  $\pi$ -pulses, applied to the atoms in reverse order, returns all population to the initial states, but a relative change of sign has been accumulated between the state  $|x_0\rangle$  and all other components.

inversion about mean

$$c_x \rightarrow \frac{1}{N} \sum_{x'} c_{x'} - (c_x - c_{x'})$$



**Figure 3.** Grover inversion-about-the-mean: with the  $\Lambda$ -transition laser excitation scheme shown, the  $(|0_i\rangle + |1_i\rangle)/\sqrt{2}$  dark state in each atom is uncoupled while a  $\pi$ -pulse excitation transfers the bright superposition qubit state  $|B_i\rangle = (|0_i\rangle - |1_i\rangle)/\sqrt{2}$  to the Rydberg state. In a succession of such excitation pulses, if one atom is already excited, no further excitation occurs, and hence quantum register components with more atoms populating the bright states will only be excited once, while the state  $|\Psi_0\rangle$  with no bright state atoms will not be excited at all. A second set of  $\pi$ -pulses, applied to the atoms in reverse order, returns all population to the initial states, and a relative change of sign has been accumulated between the state  $|\Psi_0\rangle$  and all other components.

$$|\Psi\rangle \rightarrow U_G |\Psi\rangle$$

$$U_G = P - (I - P) = P - Q$$

$$P = |\Psi_0\rangle\langle\Psi_0| \quad |\Psi_0\rangle \equiv \sum_x \frac{1}{\sqrt{N}} |x\rangle$$

inversion about mean is equivalent to sign change conditioned on all bits being in  $(|0\rangle + |1\rangle)/\sqrt{2}$ .

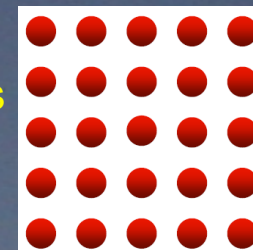
# Scaling

## Efficient Grover search with Rydberg blockade

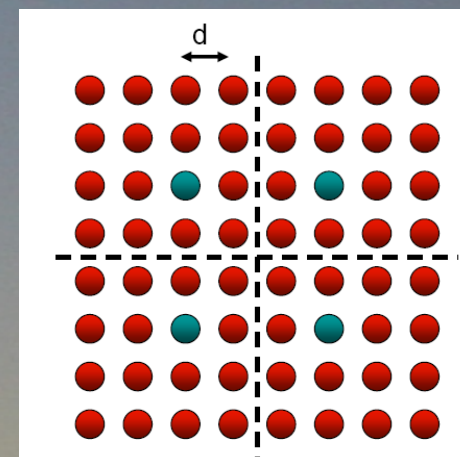
Klaus Molmer<sup>1,3</sup>, Larry Isenhower<sup>2</sup> and Mark Saffman<sup>2</sup>

- Grover iteration can be performed with 4k laser pulses, each of area  $\pi$
- Alternative, simultaneous addressing scheme requires only 8 pulses/iteration
- Best known quantum circuit: 49k-149 gates,  $k > 3$ .
- Error estimates:

$k = \log_2(N)$  qubits



sub-register architecture

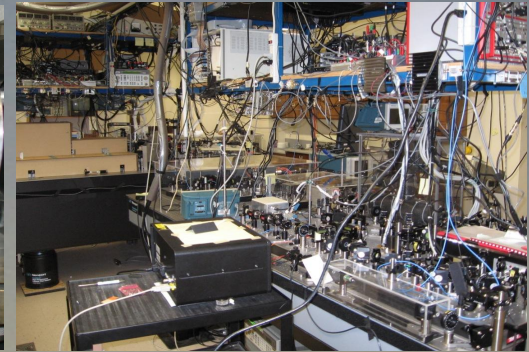
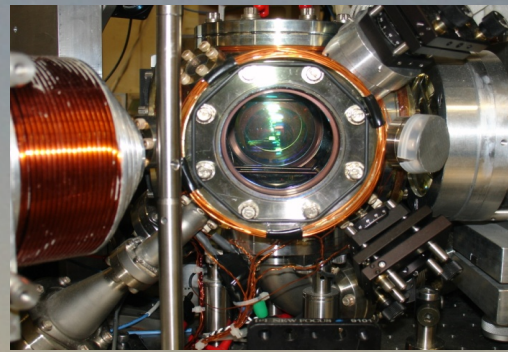
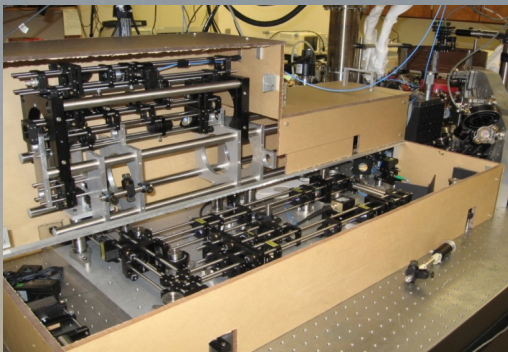
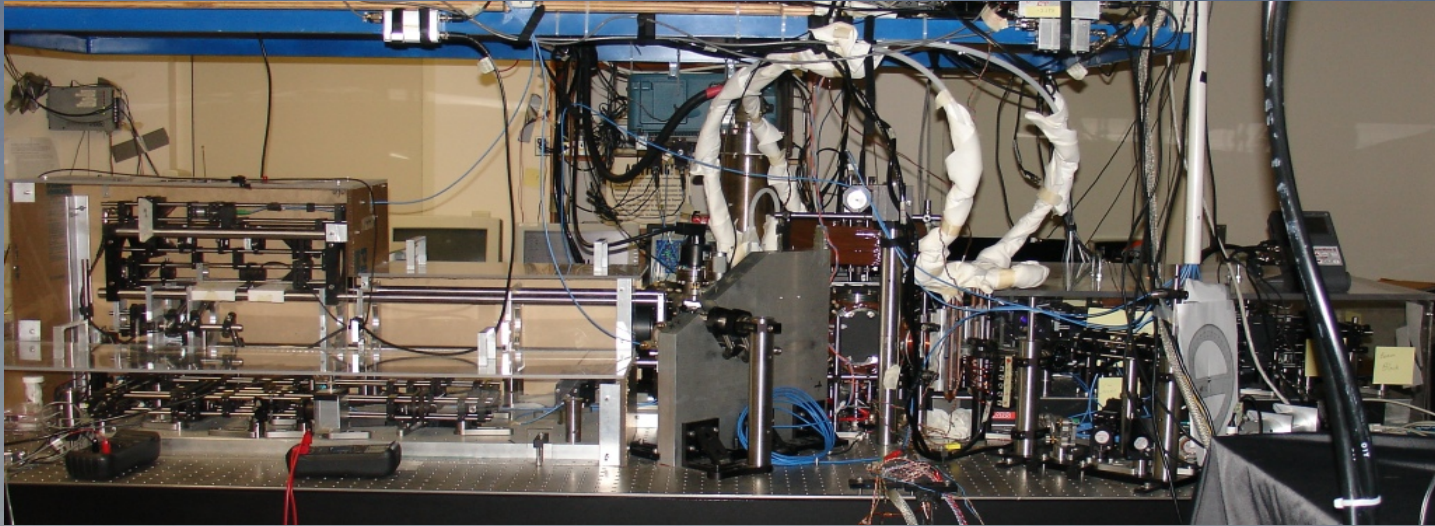


**Table 1.** Errors per Grover iteration step for the different architectural approaches described in the text.

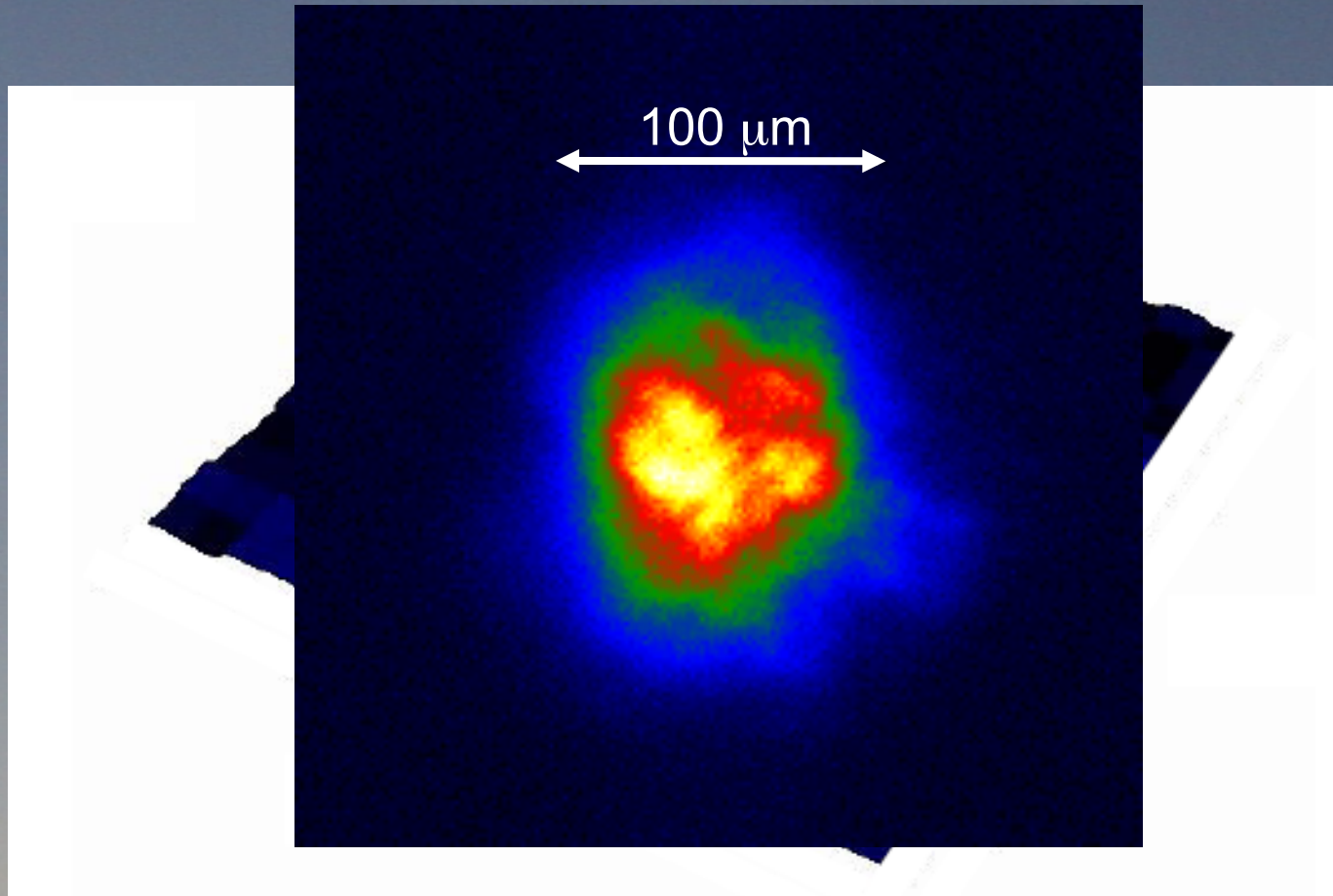
$N$	$k - 1$	$k$	Sequential addressing	Simultaneous addressing	Quadratic speedup limit $N^{-1/4}$
256	8			0.08	0.25
512		9	0.004		0.21
32 768	15			0.20	0.074
65 536		16	0.015		0.063
16 777 216	24			0.28	0.016



# Experimental Status



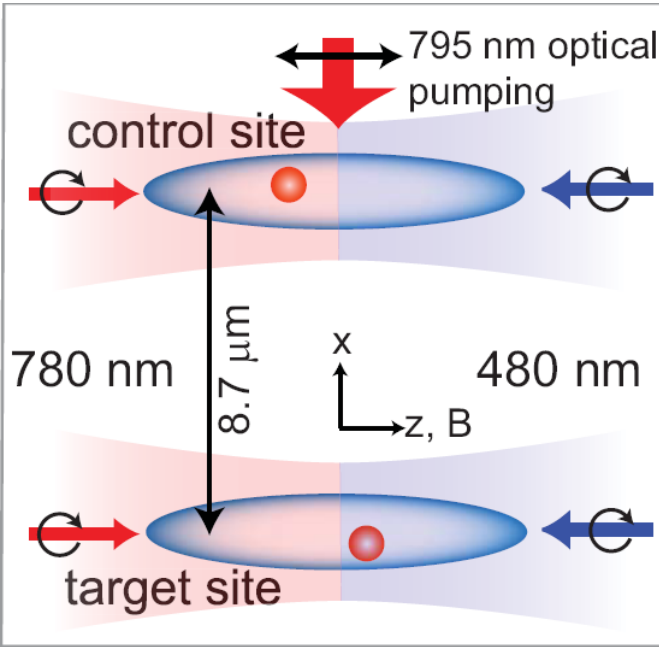
# CNOT implementation - two site experiment



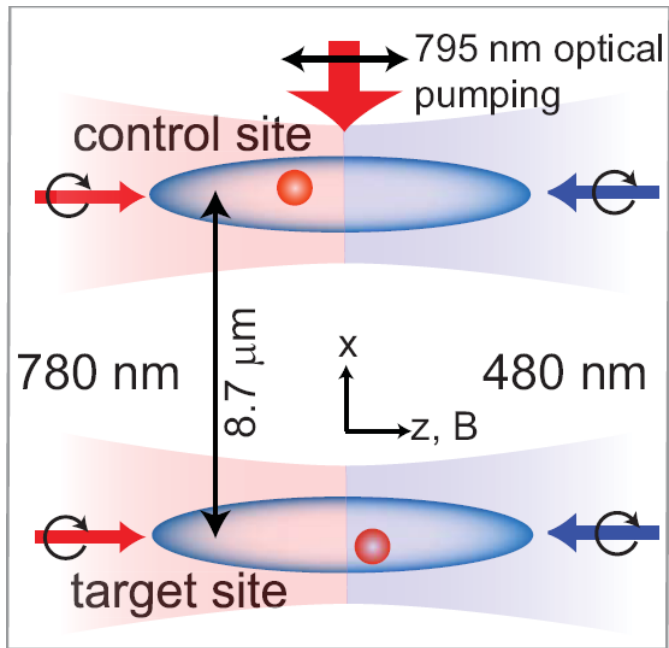
Average of 146 shots.

Probability of one atom in both sites about 10%.

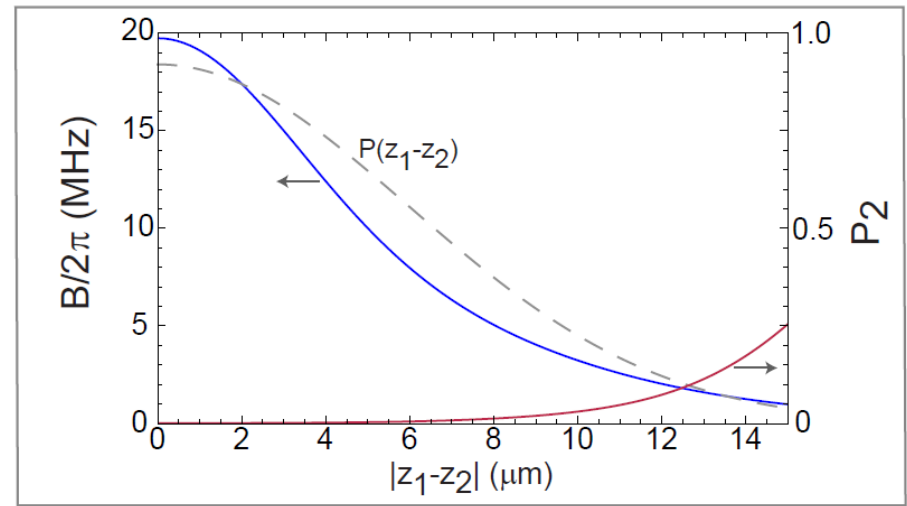
# CNOT implementation



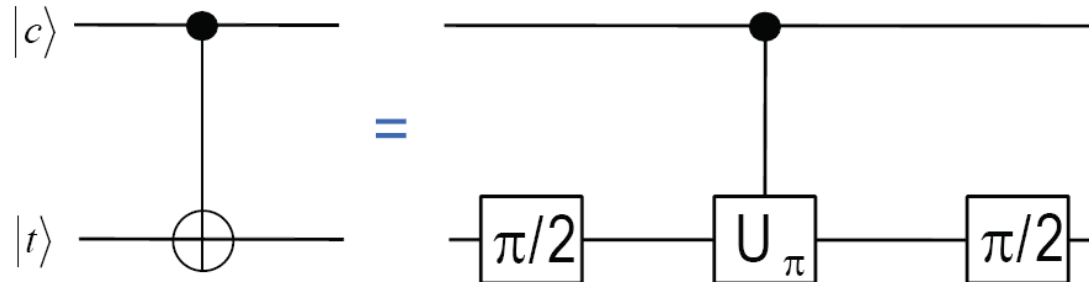
# CNOT implementation



$97d_{5/2}$

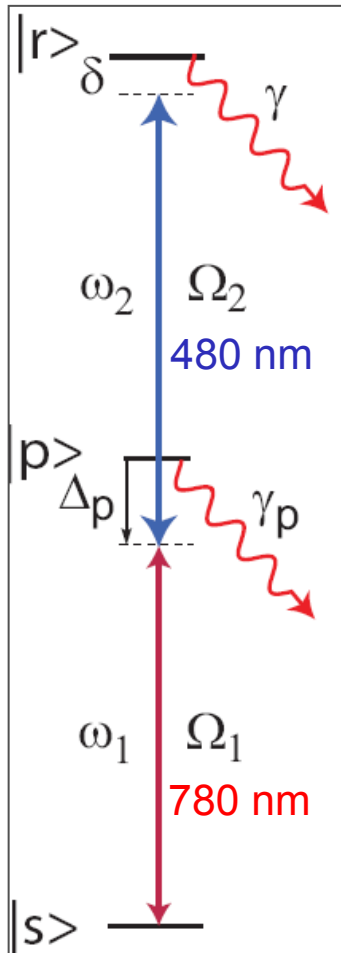


CNOT

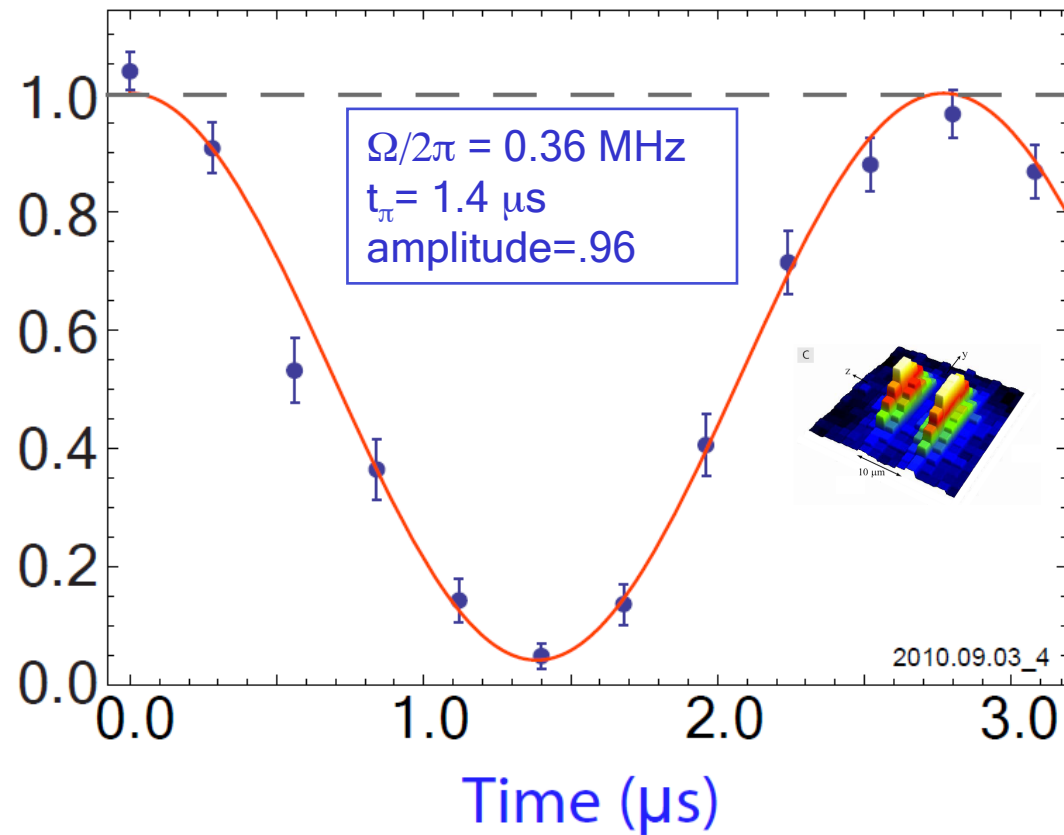


# Rydberg state Rabi oscillations

$5s_{1/2} - 5p_{3/2} - 97d_{5/2}, m=5/2$



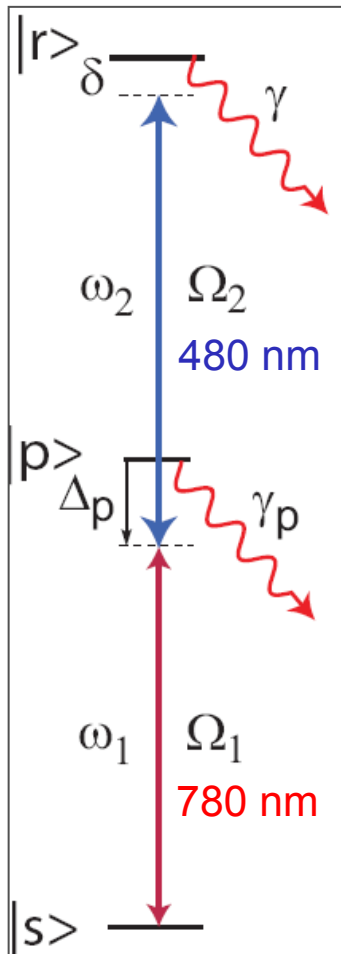
Ground State Probability



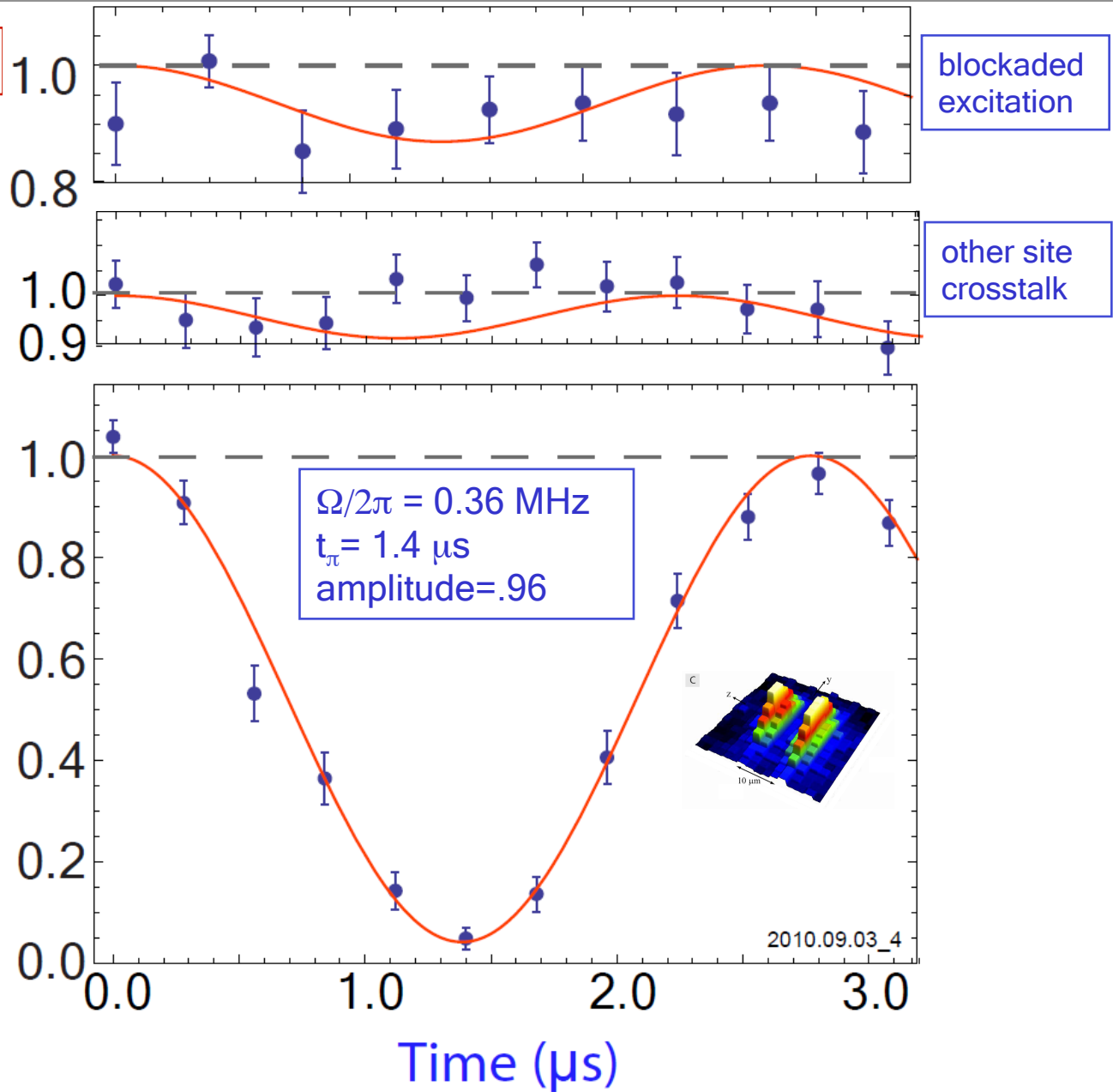
2010.09.03\_4

# Rydberg state Rabi oscillations

$5s_{1/2} - 5p_{3/2} - 97d_{5/2}, m=5/2$

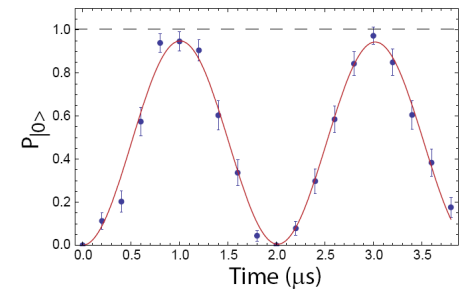
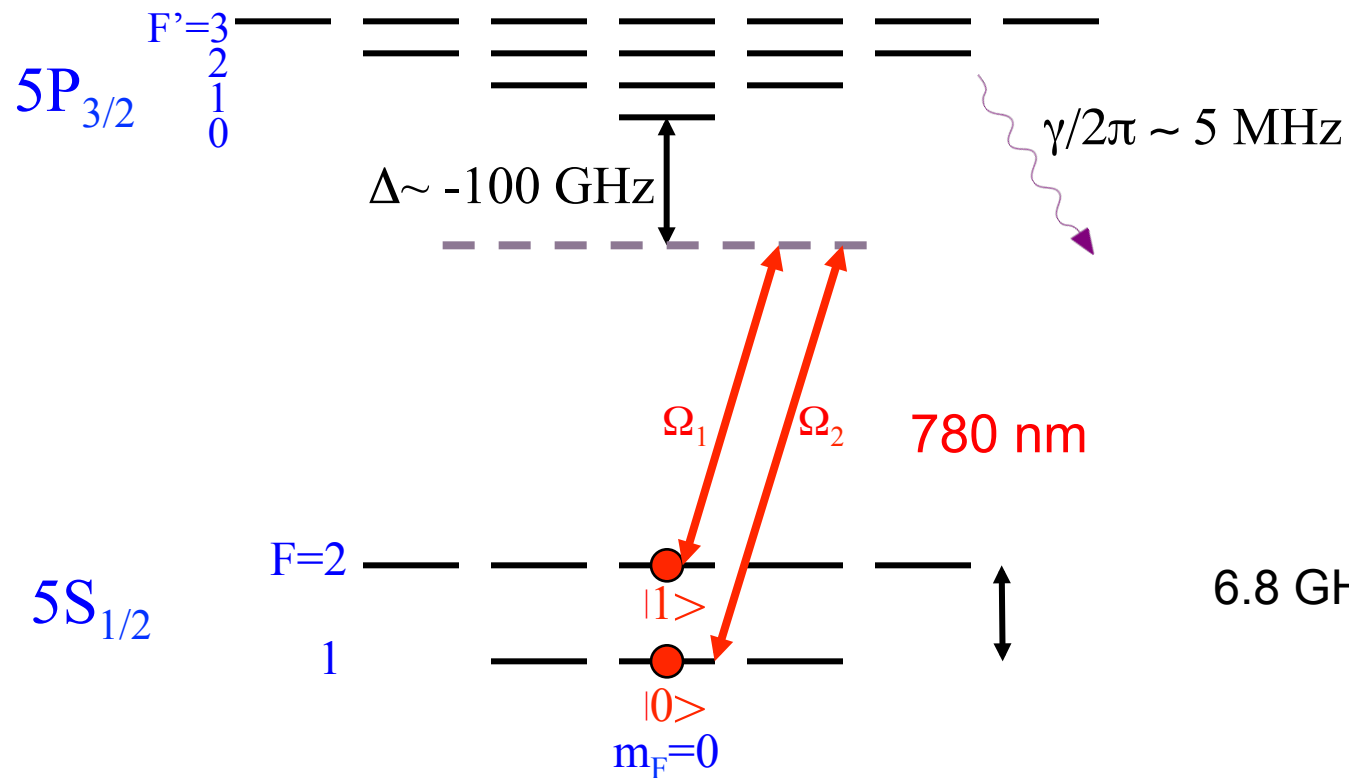


Ground State Probability

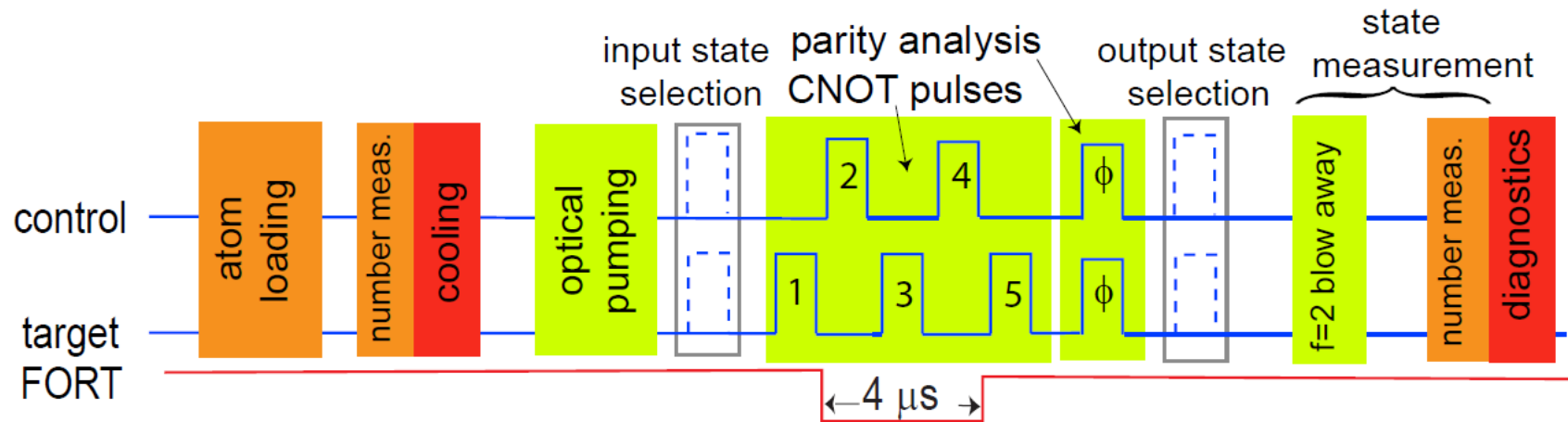
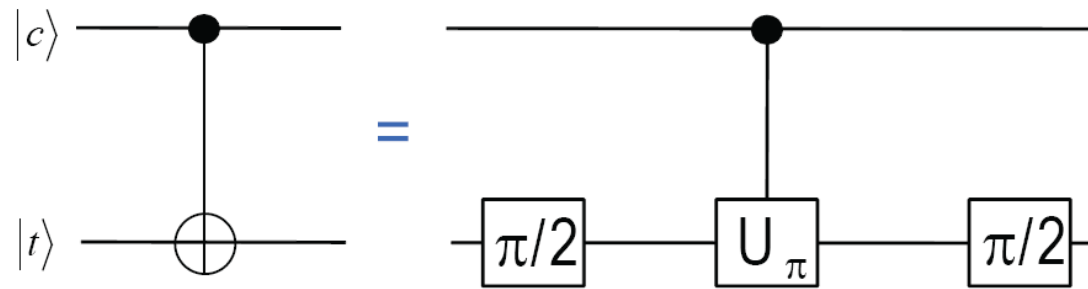


# Single qubit rotations

single qubit rotations by 2 photon stimulated Raman



# CNOT implementation





# Rydberg blockade experiment

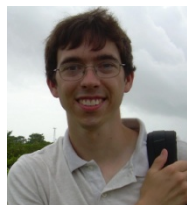
- collaboration with Thad Walker



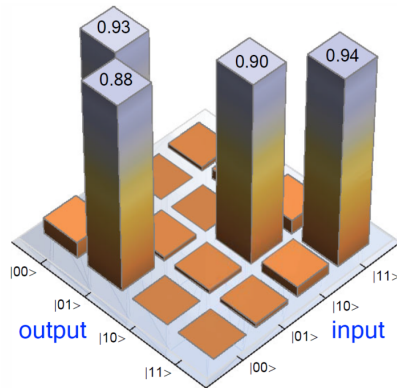
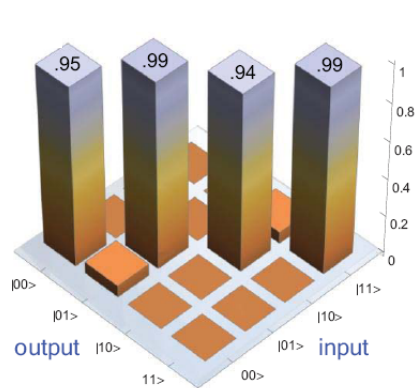
Xianli Zhang



Larry Isenhower



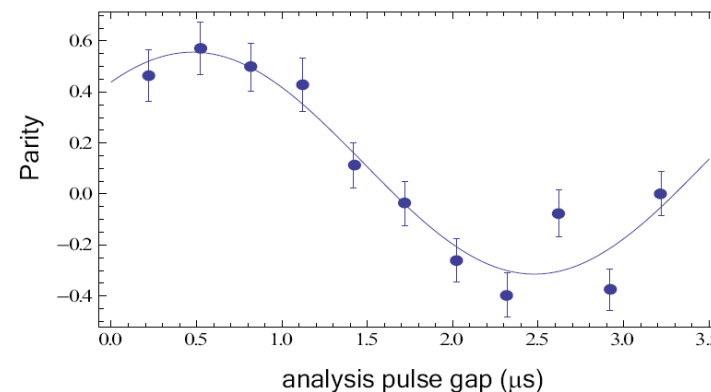
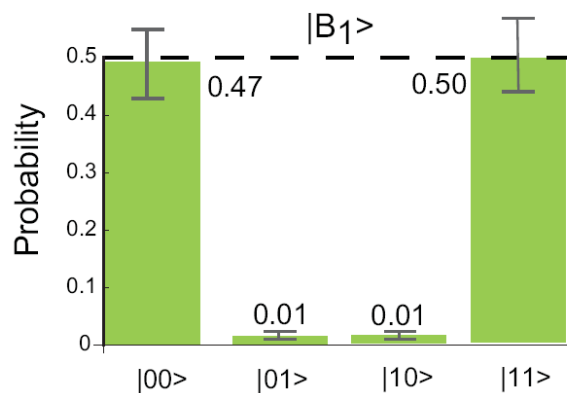
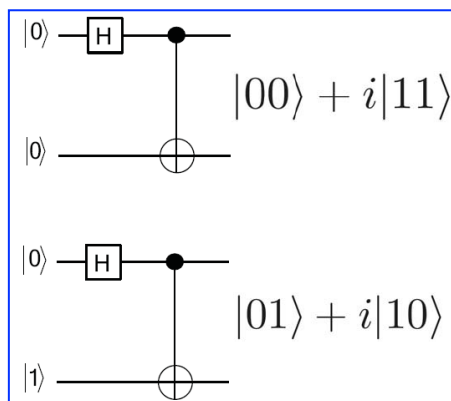
Alex Gill



**CNOT truth table**

**Fidelity: raw**  $(1/4)\text{Tr}(U_{ideal}^T U_{exp}) = .74 \pm .05$   
**loss corrected**  $(1/4)\text{Tr}(U_{ideal}^T U_{exp}) = .92 \pm .06$

## Entanglement

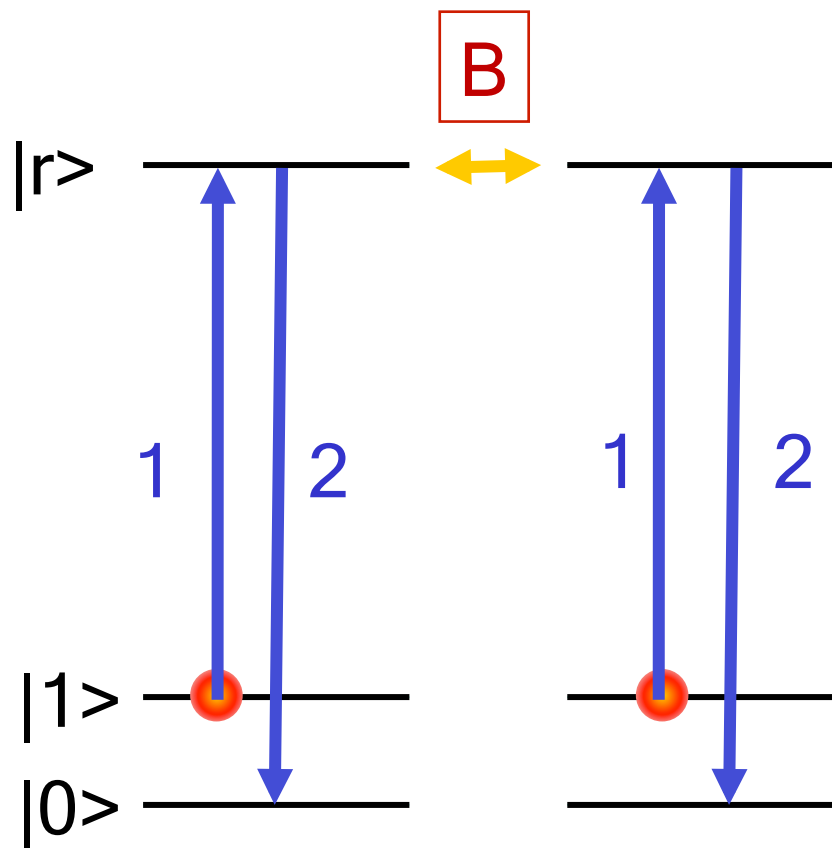


**Fidelity: raw** 0.58+/-0.04

**loss corrected** 0.71+/-0.05

Isenhower et al. PRL 2010  
 Zhang, et al. PRA 2010

# Entanglement by global addressing



$$|\psi\rangle = |11\rangle$$

Inst. d'Opt.: Wilk, et al. PRL (2010)

## Summary: Rydberg mediated quantum logic experiments

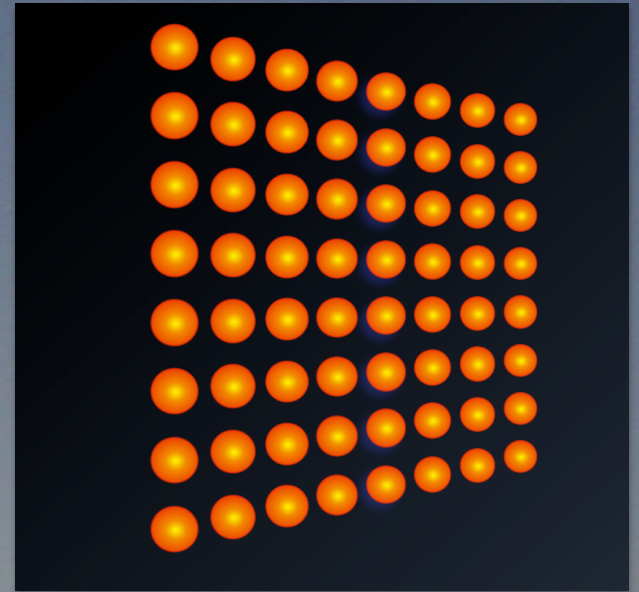
### Results with two qubits:

Single qubit operation	F ~ 0.95
CNOT truth table	F ~ 0.9
Entanglement fidelity	F ~ 0.7
Gate time	~ 5 $\mu$ s
Coherence time	~ 40 ms
Theoretical fidelity limit	~0.001
Fault tolerance threshold (architecture dependent)	.01 - .0001

# Atomic Qubit Array wishlist

Single site operations without crosstalk:

- atom loading (reloading)
- fiducial state preparation
- single and two qubit gates with high fidelity
- QND Measurements (free space QND  
Browaeys, Chapman)
- recooling
- long coherence time
- parallel operations



special requirement – trapping of both ground and Rydberg states

# Ground-Rydberg magic trapping

PHYSICAL REVIEW A **84**, 043408 (2011)

Magic-wavelength optical traps for Rydberg atoms

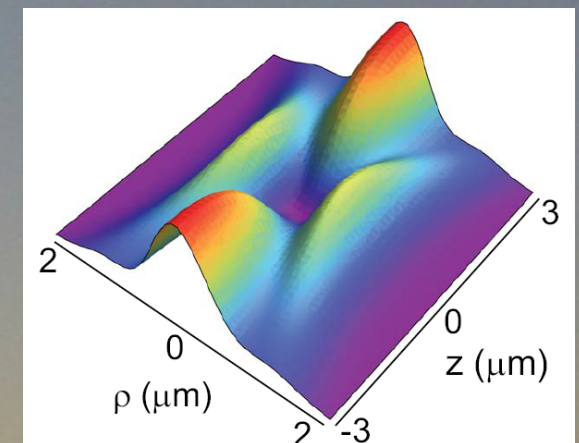
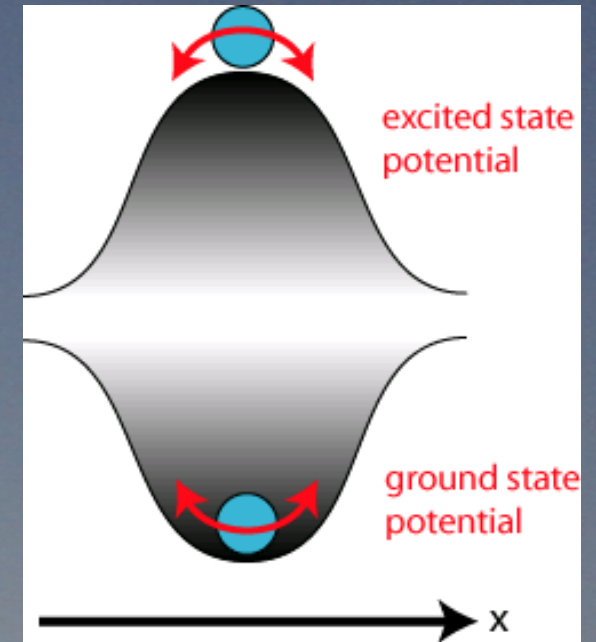
S. Zhang,<sup>1</sup> F. Robicheaux,<sup>2</sup> and M. Saffman<sup>1,\*</sup>

Rydberg state has negative polarizability  
(free electron, ponderomotive  
potential, Raitel)

Ground state has positive  
polarizability for 1 micron  
wavelength light

Choose wavelength so also ground state  
has negative polarizability

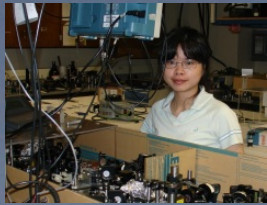
Blue detuned trap with same potential for ground and  
Rydberg states



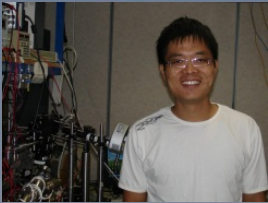
# Single atom in BoB trap

A crossed vortex bottle beam trap for single-atom qubits

G. Li, S. Zhang, L. Isenhower, K. Maller, and M. Saffman\*



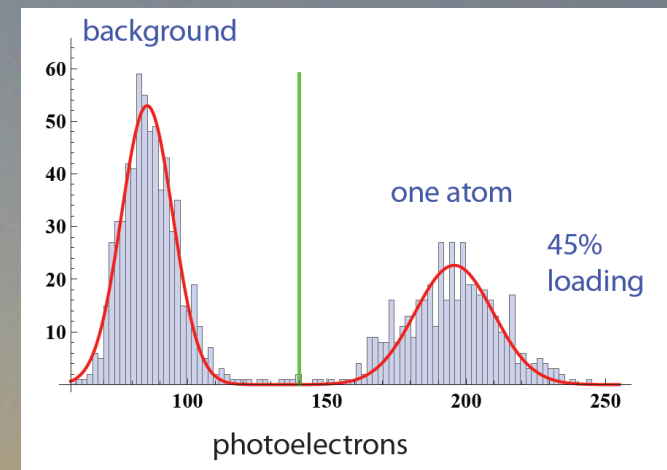
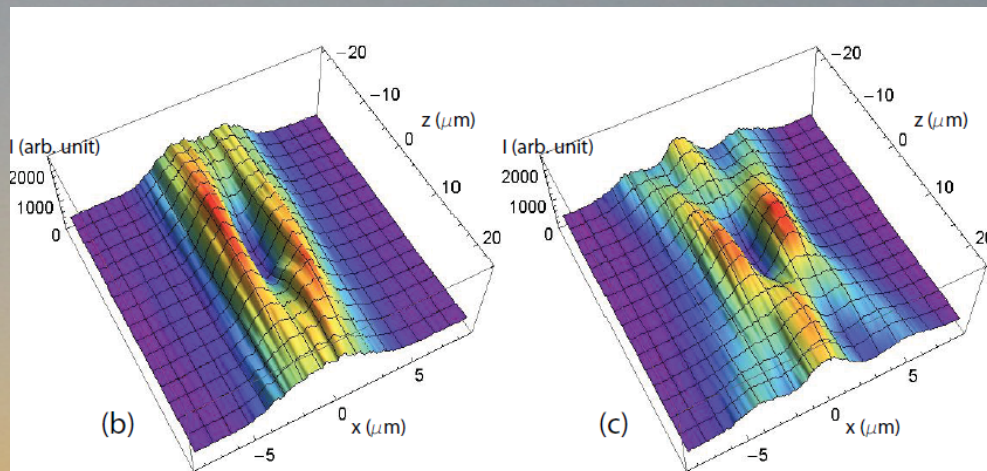
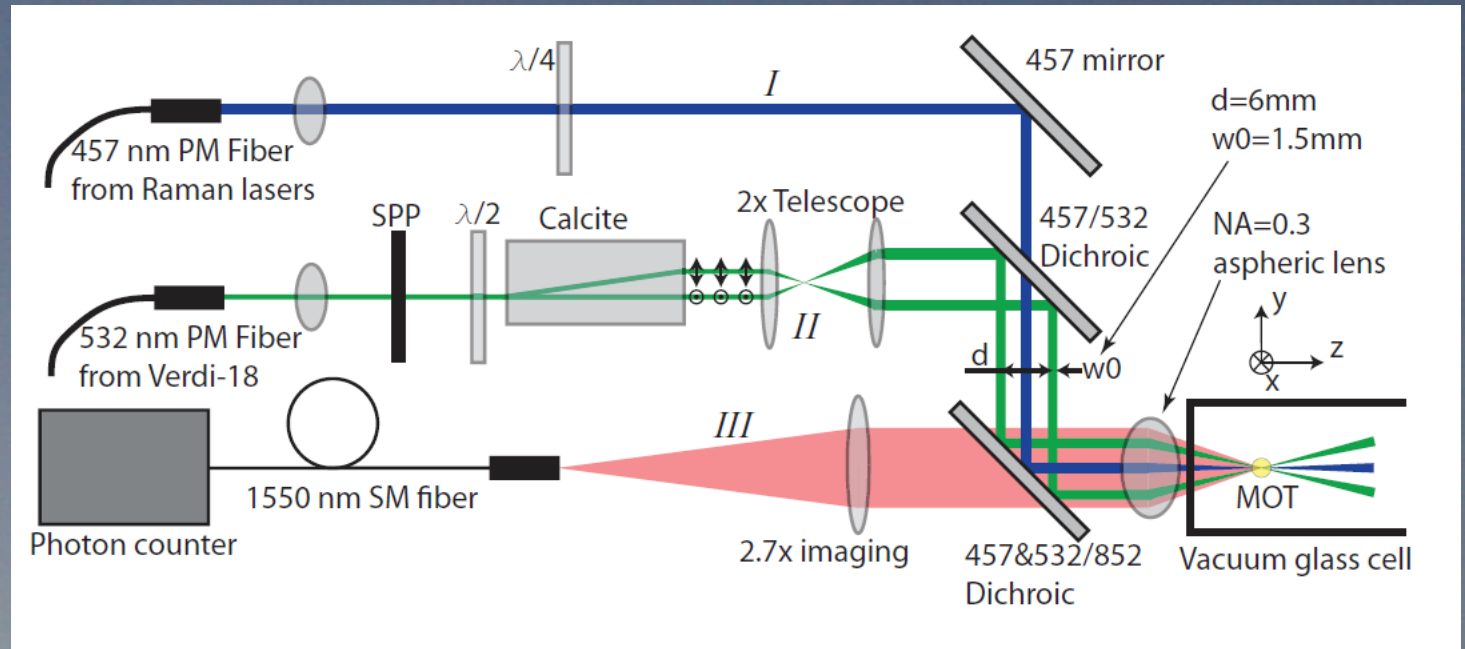
Siyuan Zhang



Gang Li

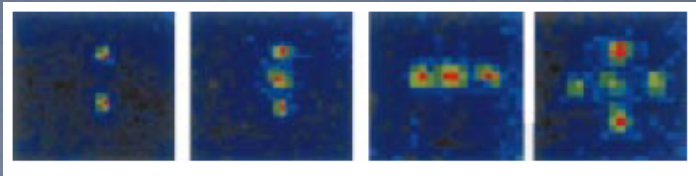


Kara Maller

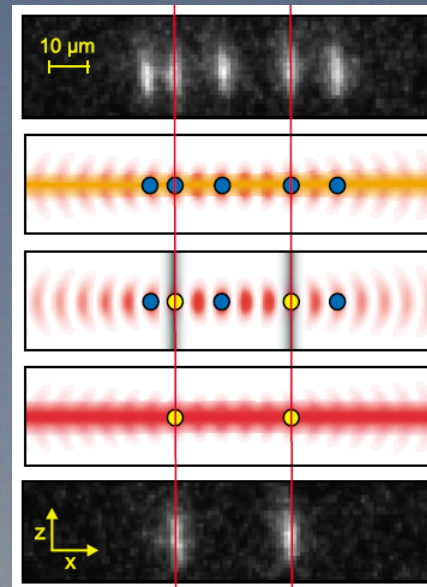


# Atomic Qubit Array

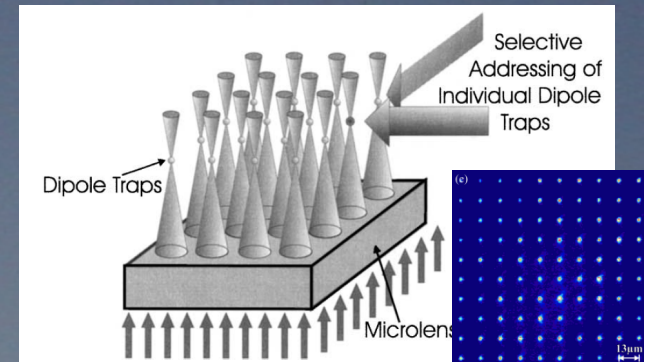
Orsay 2004



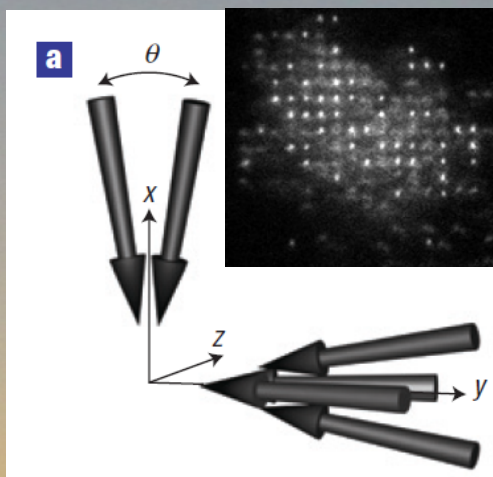
Bonn 2004



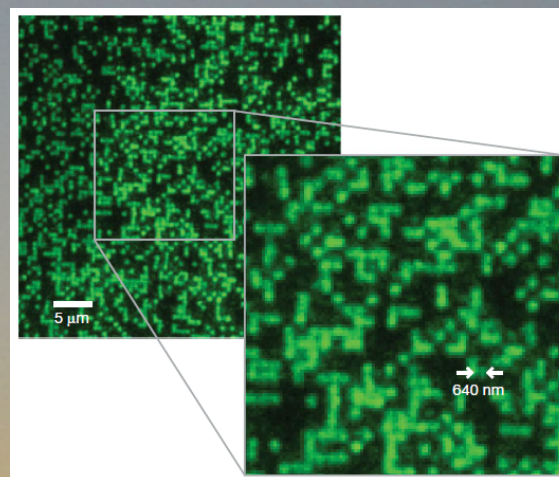
Darmstadt 2010



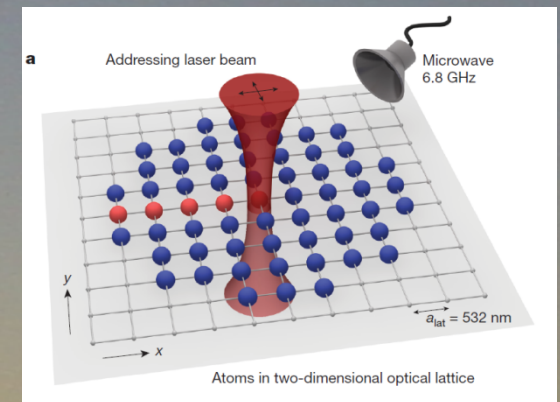
Penn State 2007



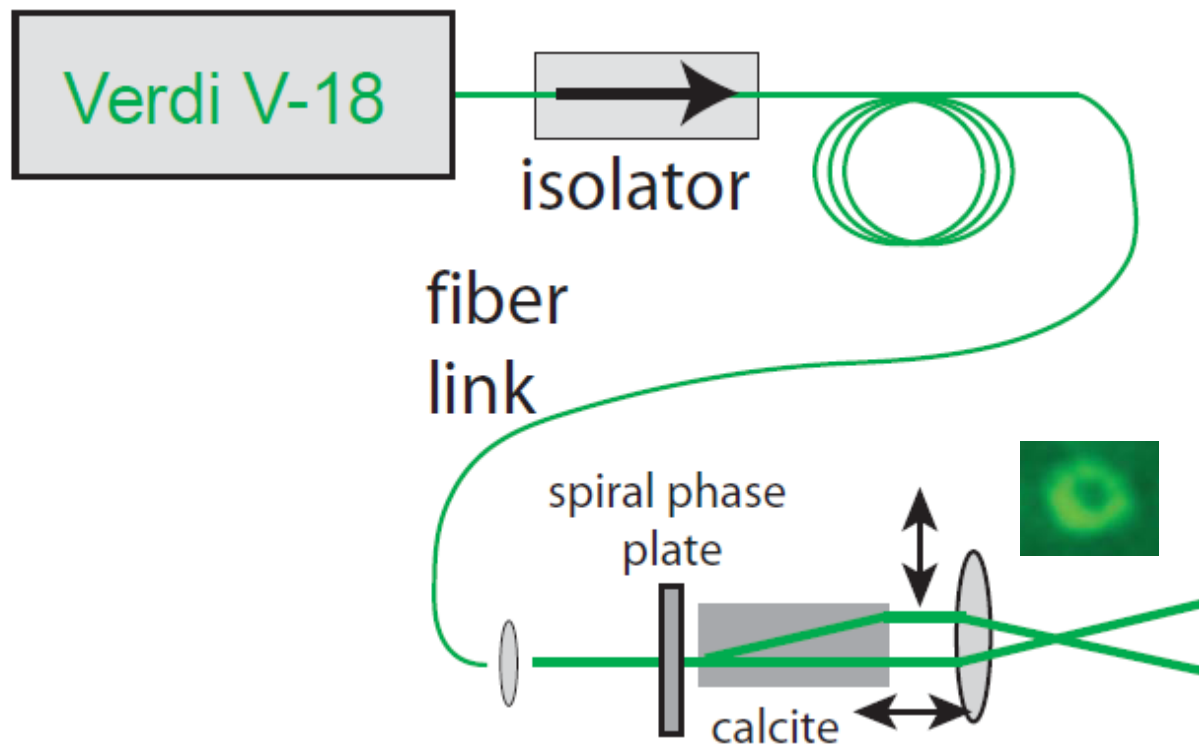
Harvard 2009



Munich 2011



# Array of bottle traps





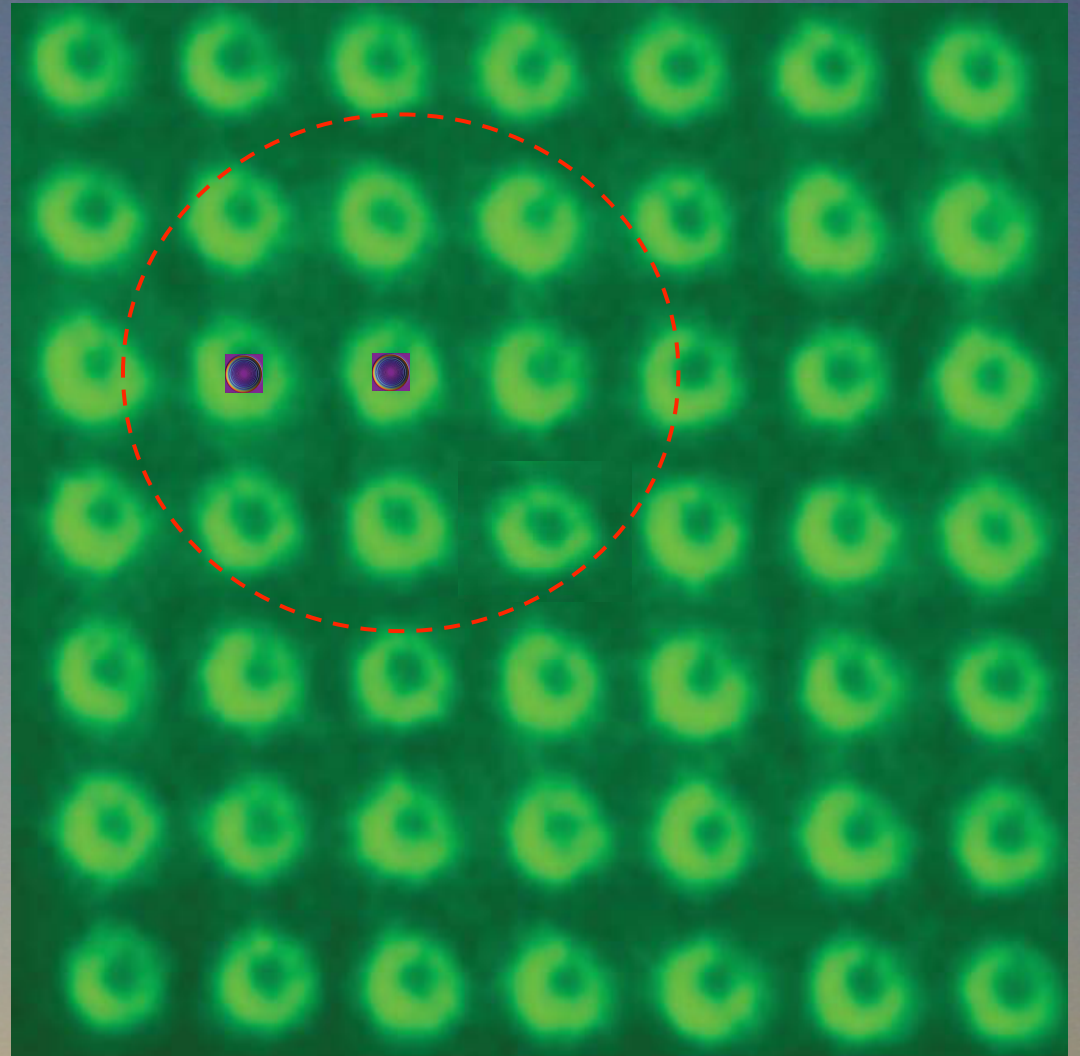
# Array of bottle traps

Two walls between  
each trap

$N=100 \rightarrow$  spacing  $\sim 8 \mu\text{m}$

Blockade distance  $\sim 30 \mu\text{m}$

9 coupled qubits



# 2D array of traps

One wall between  
each trap

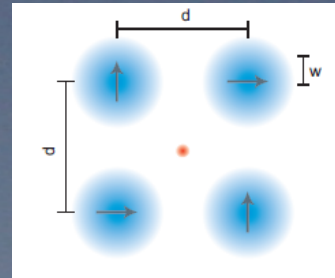
$N=100 \rightarrow$  spacing  $\sim 4 \mu\text{m}$

40 coupled qubits

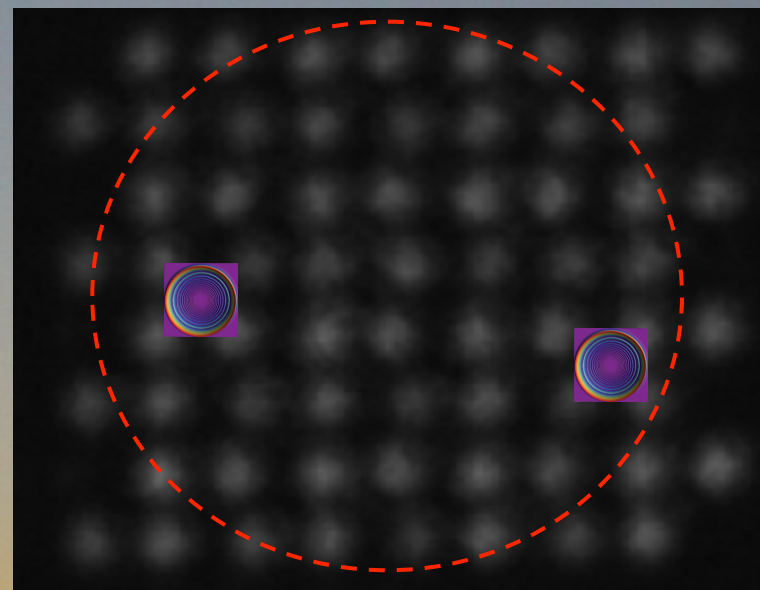
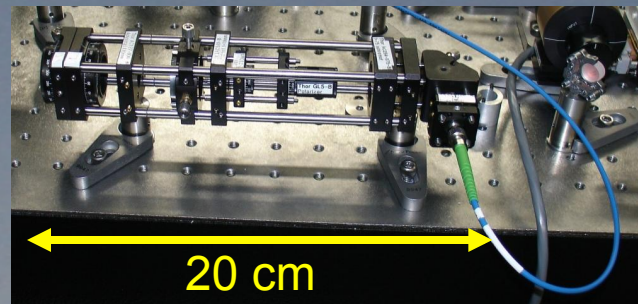
projected lattice, 42 sites

no sensitivity to phase drifts

long term stability

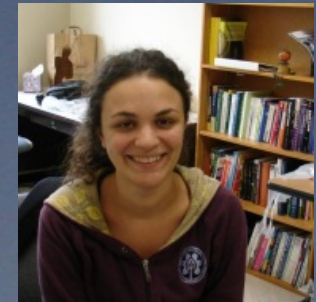
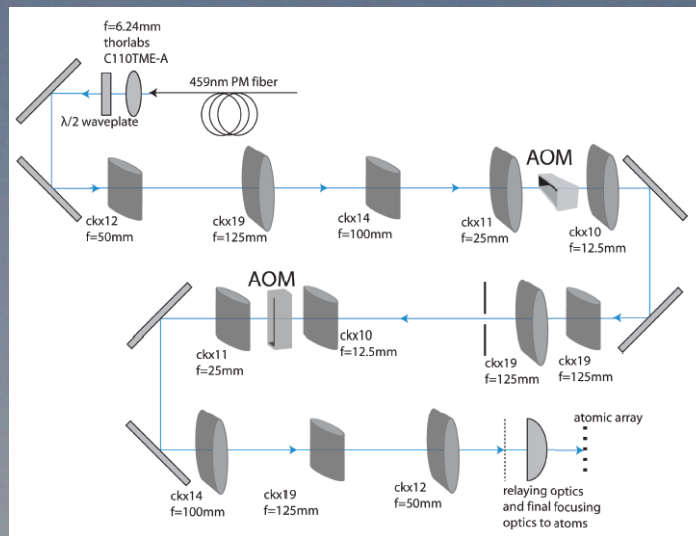


Michal  
Piotrowicz

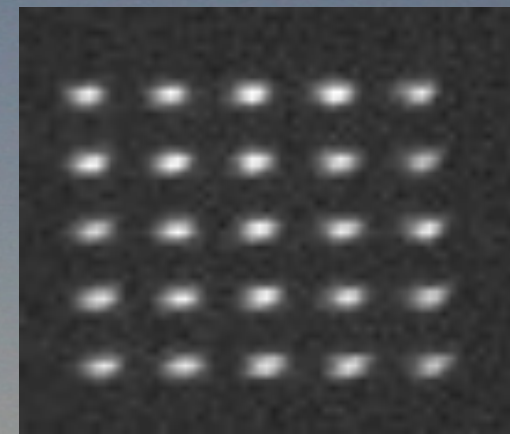
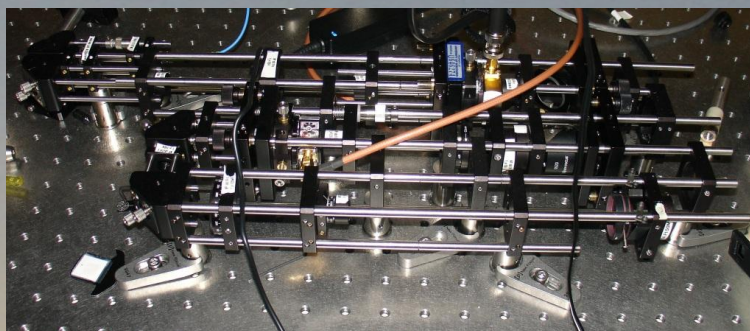


# Qubit addressing

## 2D acousto-optic beam scanners



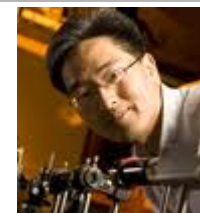
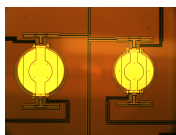
Kara Maller



25 scanned spots,  
sub  $\mu$ s addressing

# Using MEMS for atomic QIP

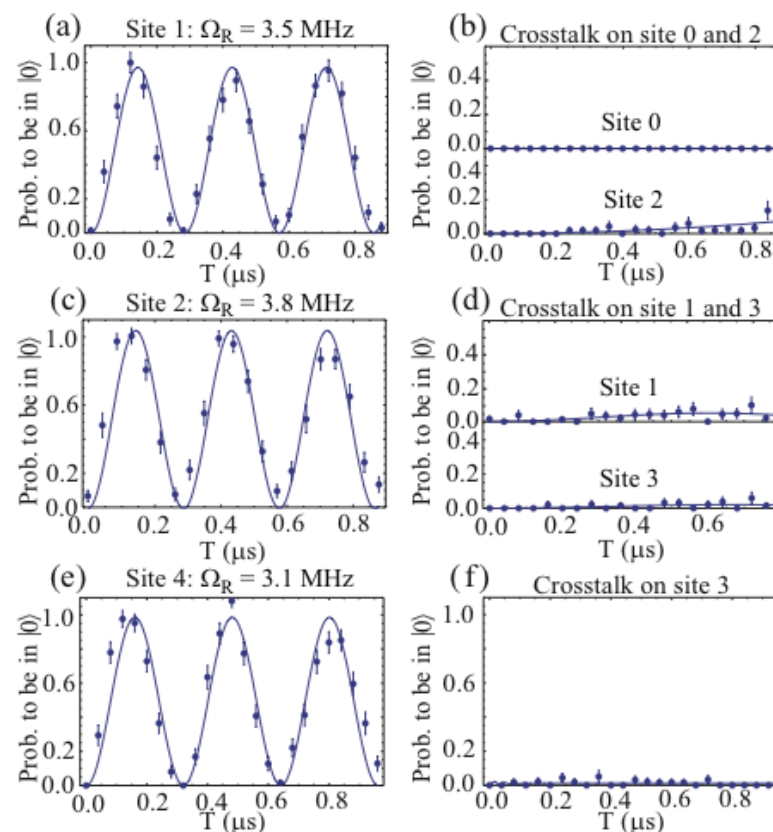
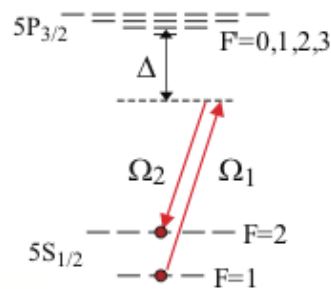
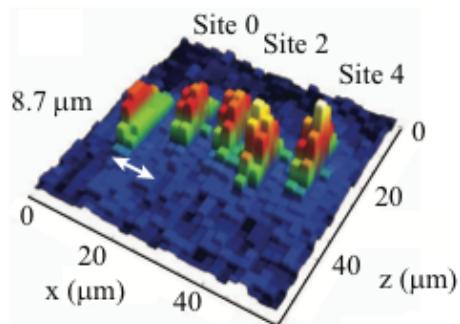
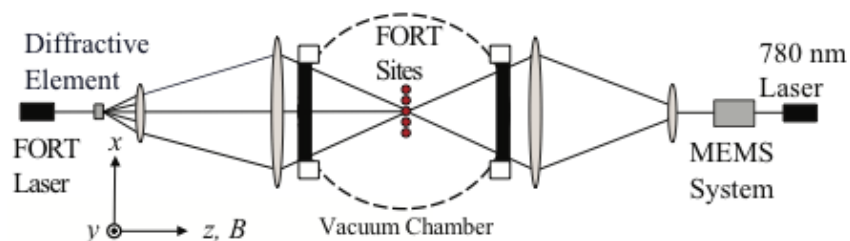
- Addressing individual Rb atoms



Jungsang Kim



Felix Lu



(U. Wisconsin, Duke, AQT)

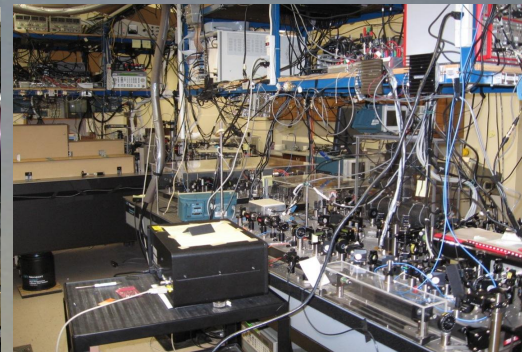
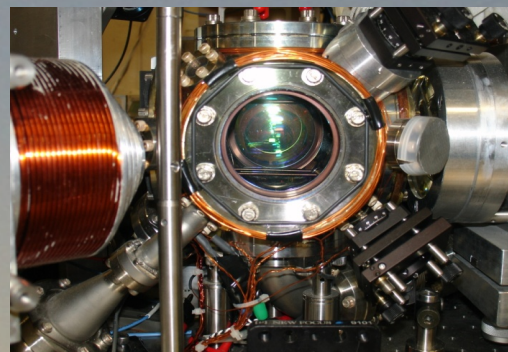
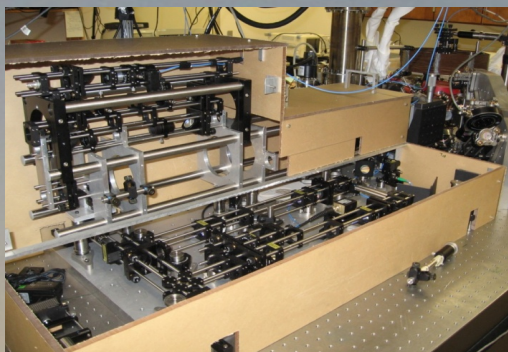
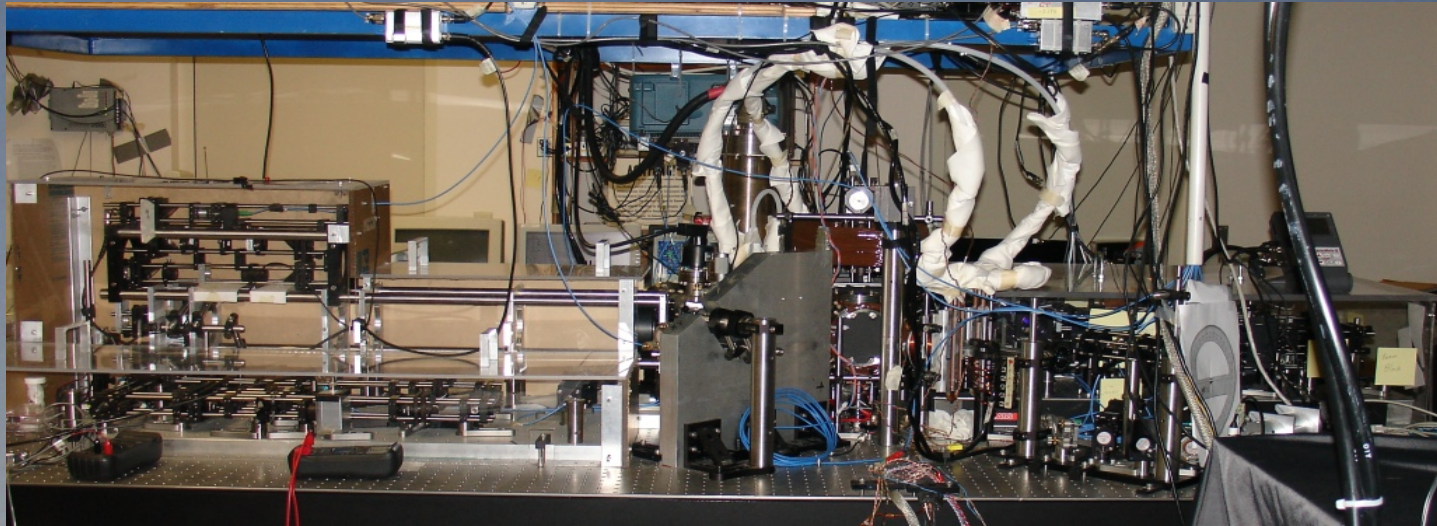
C. Knoernschild et al., Appl. Phys. Lett. (2010)

# Footprint: Rb experiment 2001-2011

3m

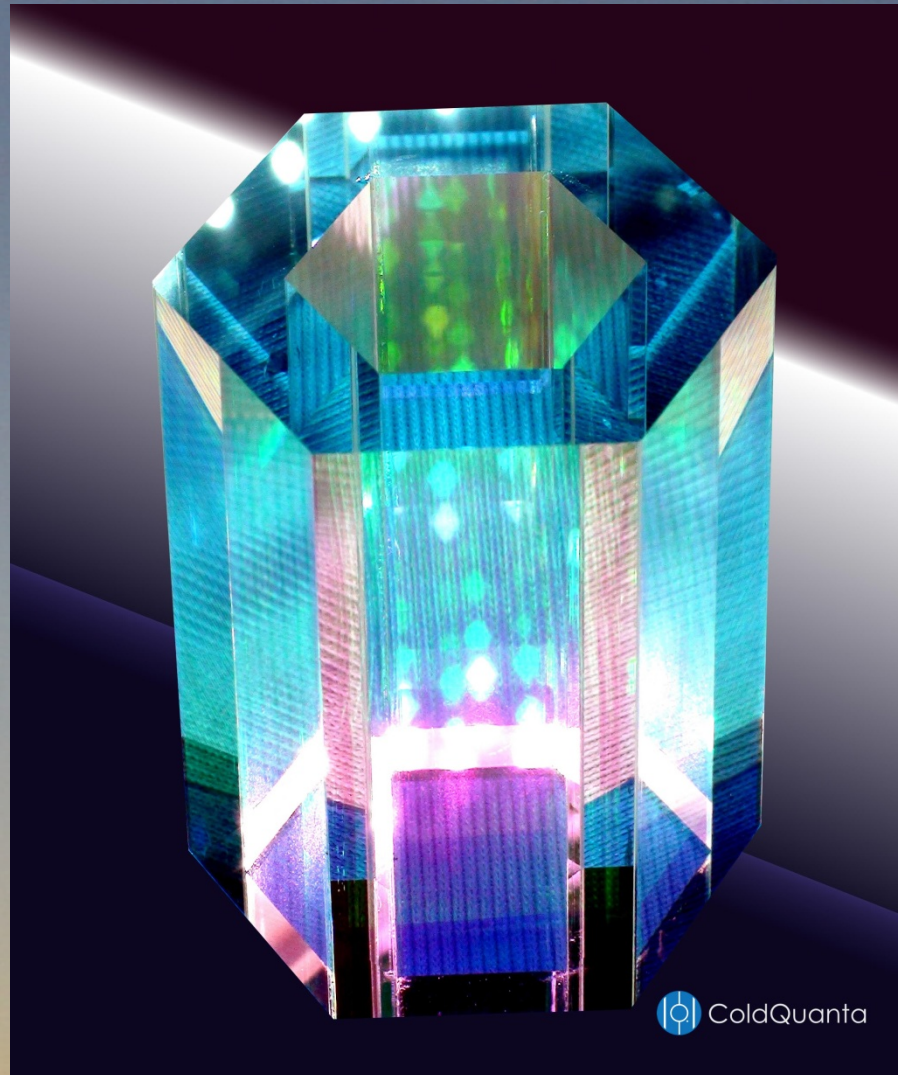
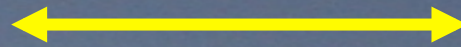


Rb



# Next generation

3.5 cm



Dana Anderson



ColdQuanta

SRI International  
SARNOFF



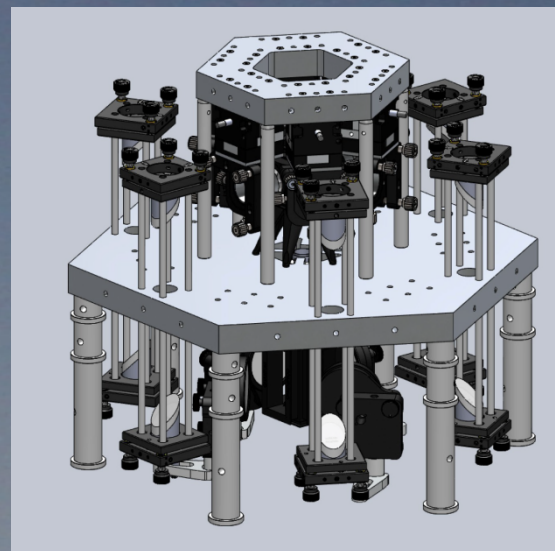
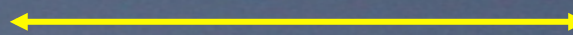
# Next generation

small pyrex cells

6 sides  $\Omega/4\pi=40\%$



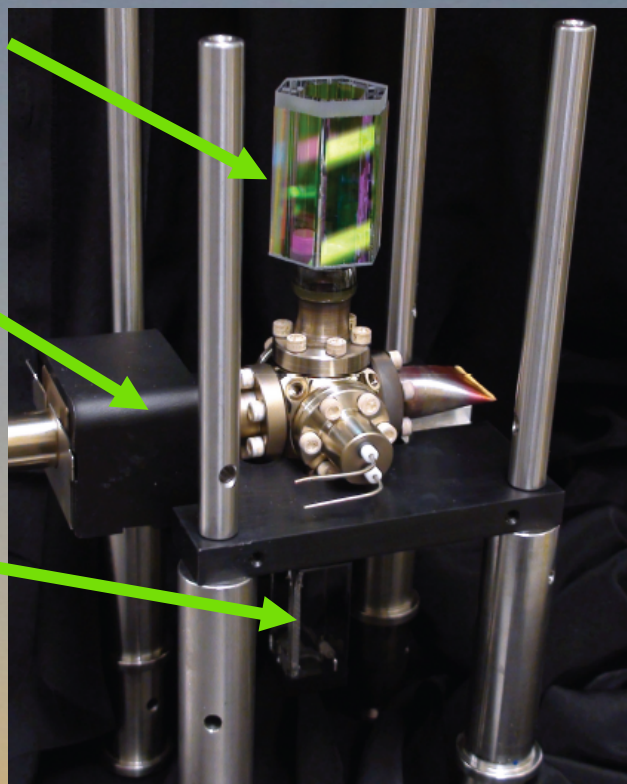
addressing optics,  $\sim 25$  cm



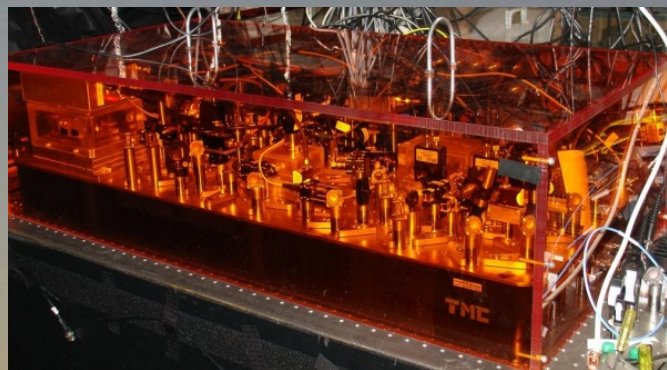
science cell

ion pump

2D MOT source



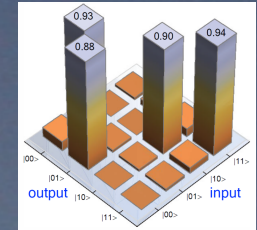
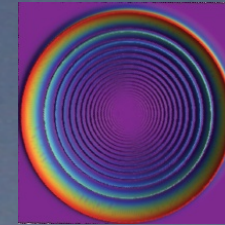
new lasers, referenced to frequency comb



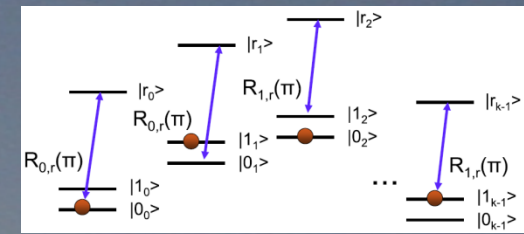
# Quantum Information with Rydberg atoms



Neutral atoms are excellent candidates for scalable QIP



Multi-bit Rydberg interactions for efficient Implementation of algorithms



New technology for multi-qubit experiments

