

A Magnetic Reconnection Mechanism for Ion Acceleration and Abundance Enhancement in Impulsive Flares

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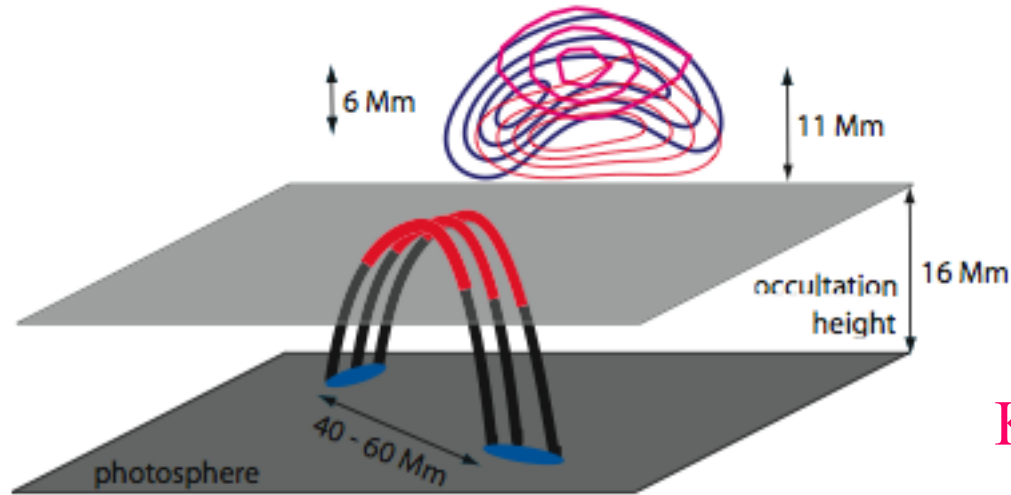
Energetic particle observations in impulsive flares

- Flare and coronal observations

- In solar flares energetic electrons up to MeVs and ions up to GeVs have been measured
 - A significant fraction of the released magnetic energy appears in the form of energetic electrons and ions (Lin and Hudson '76, Emslie et al '05, Krucker et al '10)
 - A large number of electrons undergo acceleration – the “numbers problem”
 - Correlation between > 300keV energetic electrons and > 30 MeV ions (Shih et al 2008) ⇒ common acceleration mechanism
- In impulsive flares see enhancements of high M/Q ions as well as ^3He (Mason '07)
- In the extended corona $T_{\perp} > T_{\parallel}$ (Kohl '97)
 - Minority ion temperature more than mass proportional

$$T_i / T_p \geq m_i / m_p$$

RHESSI occulted flare observations



30-50keV

17GHz

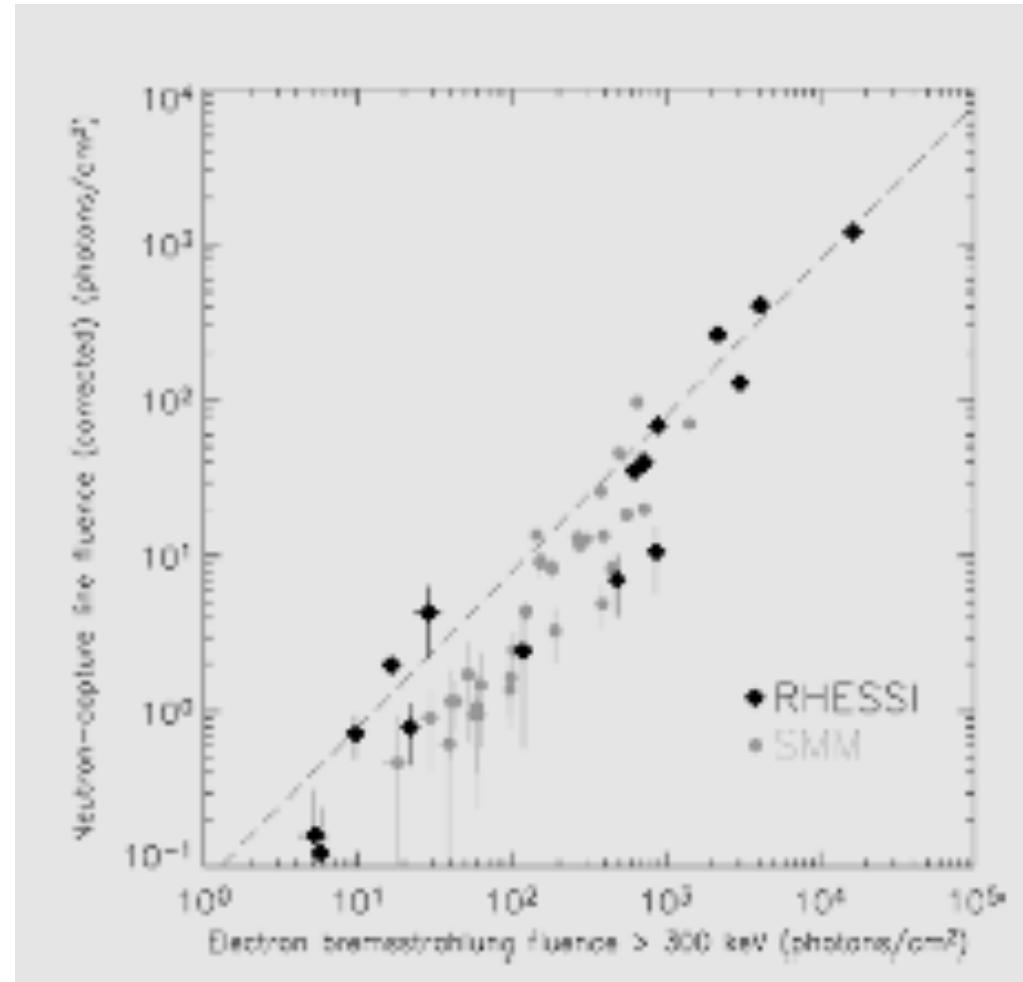
Krucker et al 2010

- Observations of a December 31, 2007, occulted flare
 - All electrons in the flaring region are part of the energetic component (10keV to several MeV)
 - The pressure of the energetic electrons approaches that of the magnetic field
 - Conversion of magnetic energy to energetic particles is very efficient
 - Remarkable!

Energetic electron and ion correlation

- $> 300\text{keV}$ x-ray fluence (electrons) correlated with 2.23 MeV neutron capture line ($> 30\text{ MeV}$ protons)
- Acceleration mechanisms of electrons and protons linked?

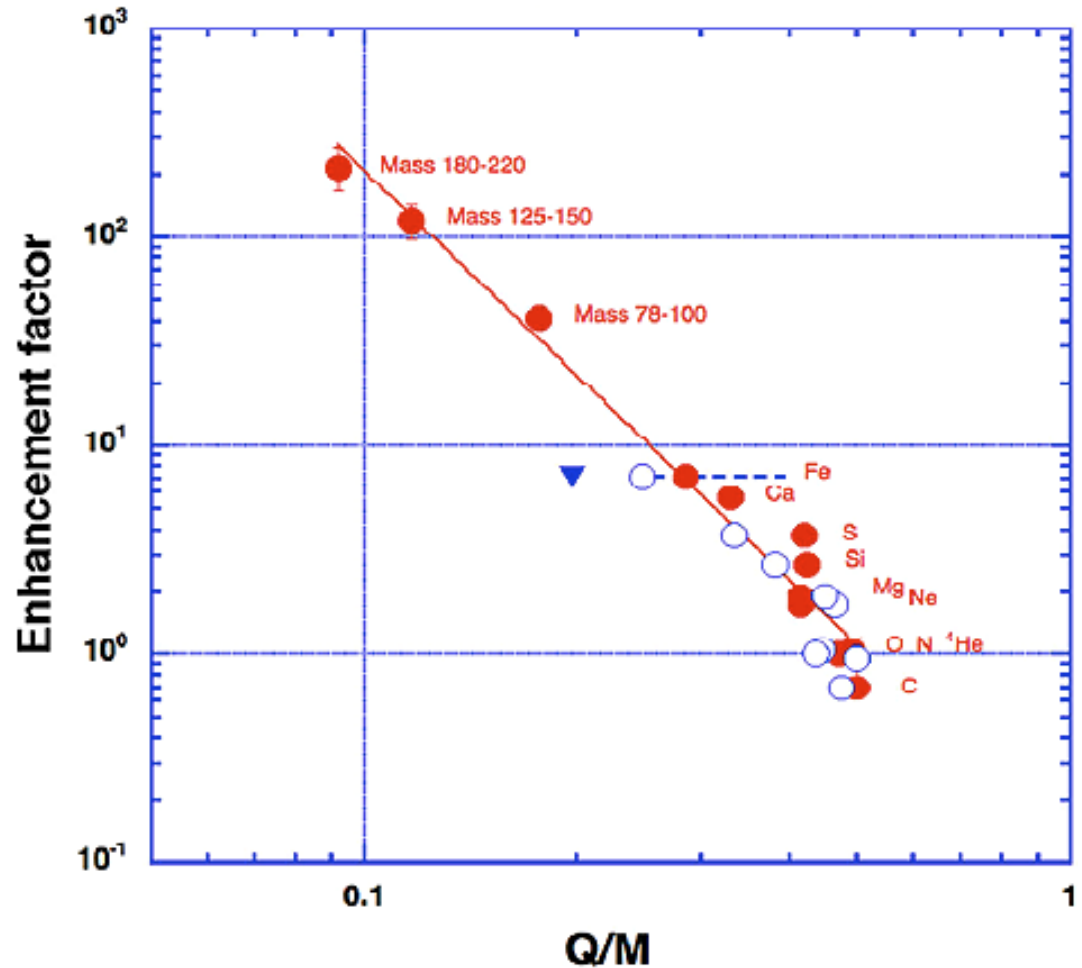
Shih et al 2008



Impulsive flare energetic ion abundance enhancement

- During impulsive flares see heavy ion abundances enhanced over coronal values
- Enhancement linked to Q/M

$$\propto \left(\frac{Q}{M} \right)^{-3.26}$$

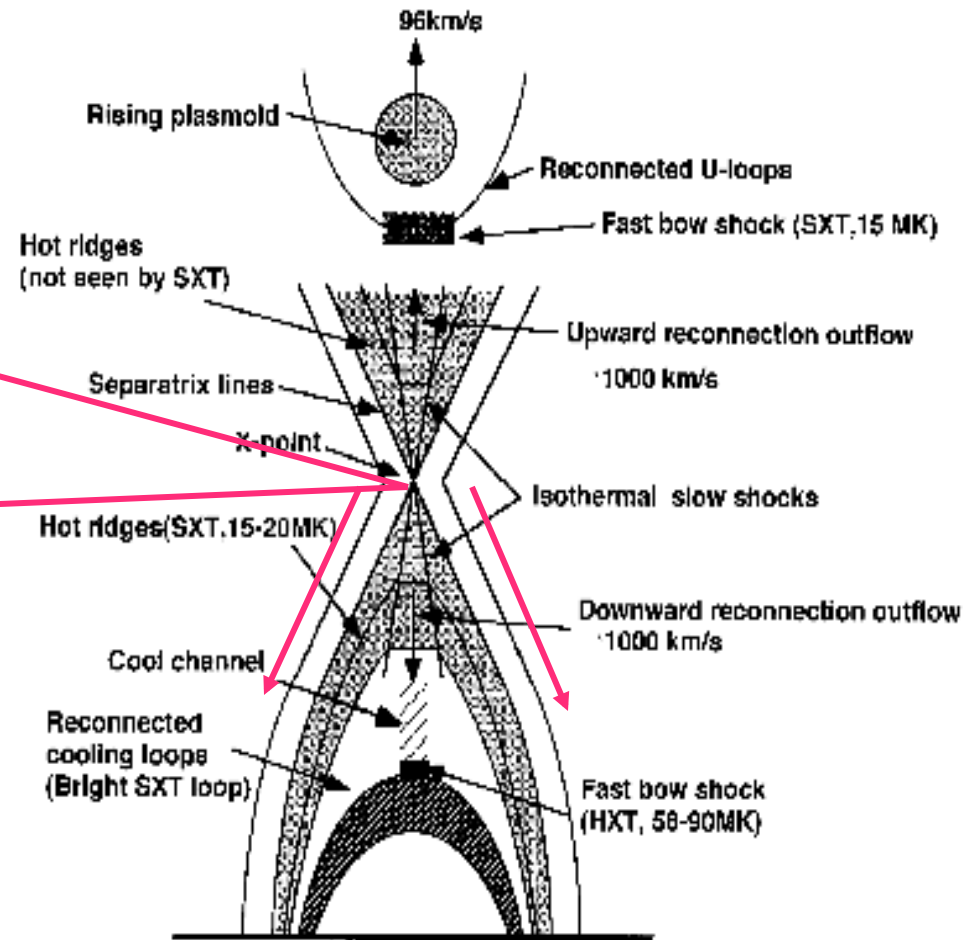
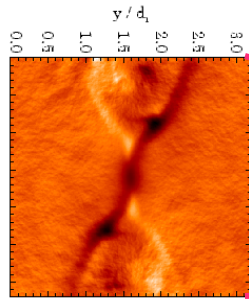


Mason, 2007

Single x-line model: the sun

- Can parallel electric fields in a single x-line produce the large number of electrons seen in flares?

- Around 10^{37} electrons/s
- Downflow currents in a single x-line would be enormous
 - Producing 10^9 G fields for $L \sim 10^4$ km
- Parallel electric fields are shorted out except around the x-line



- Magnetic energy is not released at the x-line but downstream as the reconnected fields relax their stress

- The x-line is not where energy is released
- The x-line region has negligible volume

- Can't explain the large number of energetic electrons

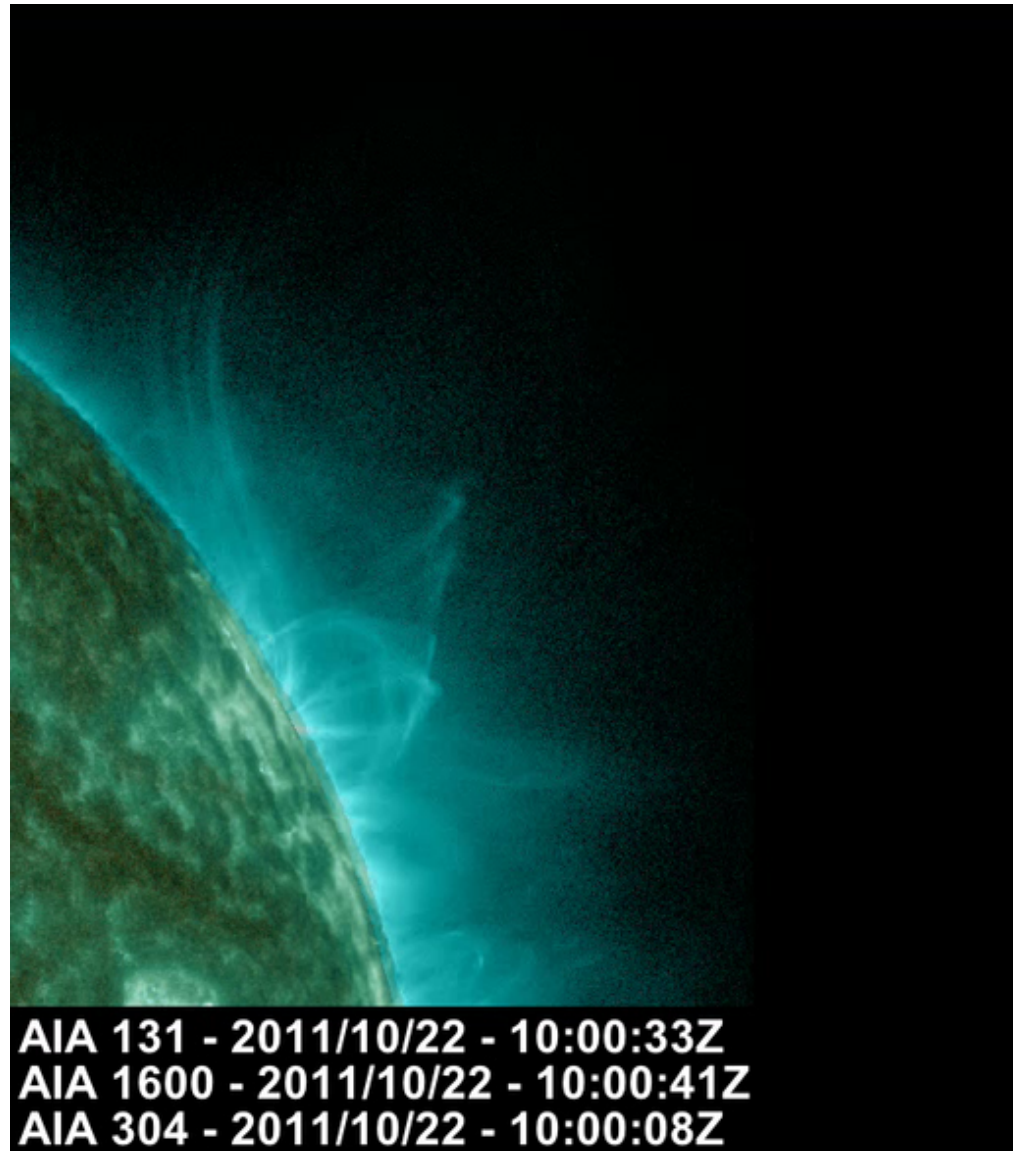
Must abandon single x-line model!

Tsuneda 1997

SDO/AIA flare observations

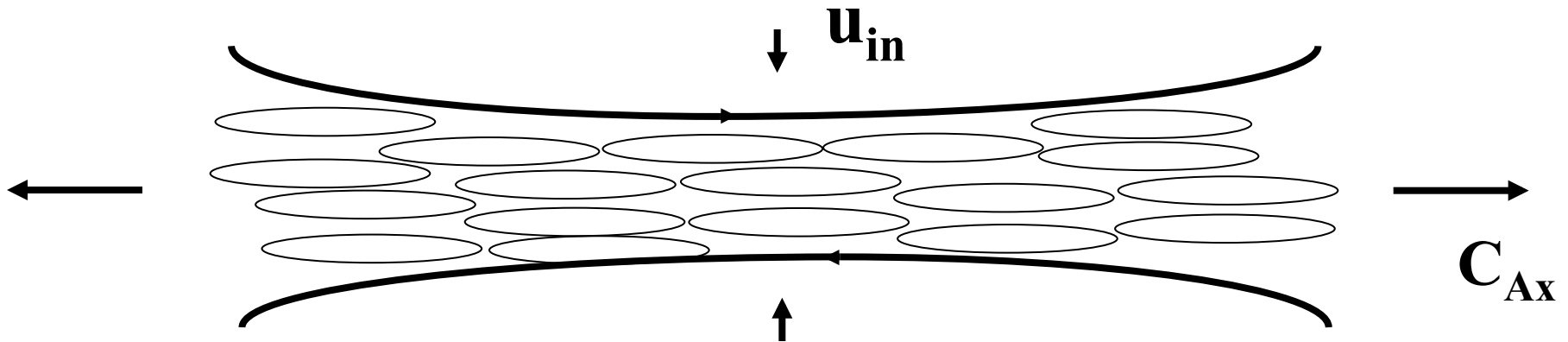
- Super Arcade Downflows (SADs) are interpreted as signatures of magnetic islands from an overlying reconnection site (Sheeley et al 2004)
- Such SAD events are now considered typical and not anomalies.
- **Must abandon the classical single x-line picture!!**

Savage et al 2012



A multi-island acceleration model

- A single x-line model can not explain the high fraction of energy going into the very energetic electrons and ions in flares
 - Must abandon the classical single x-line picture!!
- Current layers spawn multiple magnetic islands in reconnection with a guide field (Drake et al '06, Onofri et al '06, Daughton et al '11)
 - ⇒ islands have a huge range of scales
 - ⇒ SADs in flares



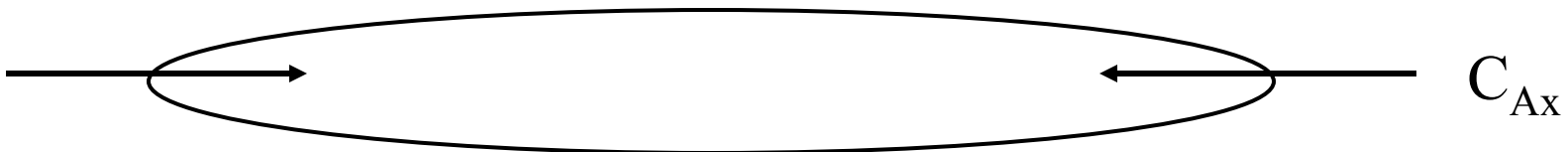
Particle acceleration in multi-island reconnection

- How are electrons and ions accelerated in a multi-island environment?
 - Fermi reflection in contracting magnetic islands (Kliem 94, Drake et al 2006, Drake et al 2010)

$$\frac{d\varepsilon_{\parallel}}{dt} \sim 2\varepsilon_{\parallel} \frac{c_A}{L_x}$$

- Rate of energy gain independent of particle mass

⇒ linkage between particle energy gain and magnetic energy release



Two-stage mechanism for ion heating in flares

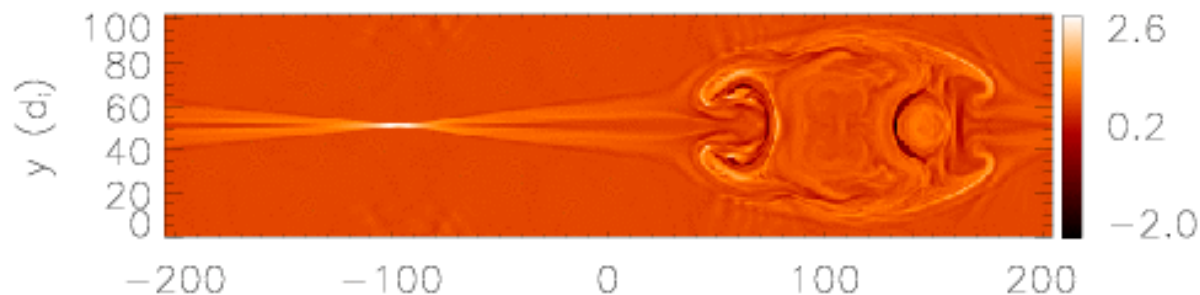
- Electrons are directly accelerated through the contracting island mechanism even in the low β corona.
- In low β conditions ions are too slow to bounce
 - Need a seed ion heating mechanism
- Stage 1: pickup of ions on entry into reconnection exhausts
 - Ions gain an effective thermal velocity comparable to the exhaust velocity

$$\Rightarrow \beta_x \sim 8\pi nT/B_x^2 \sim 1$$

- Dominantly perpendicular heating in the guide field case
- Stage 2: Fermi reflection in contracting and merging magnetic islands

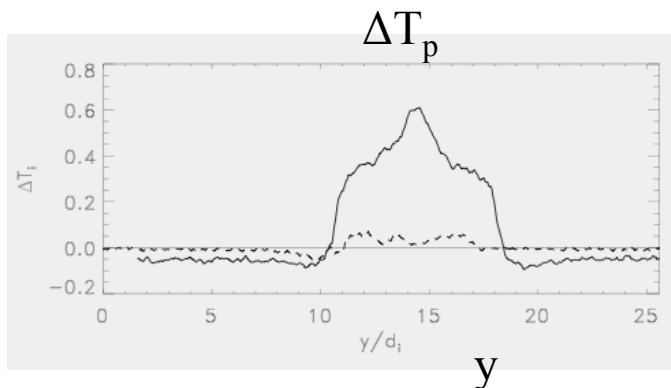
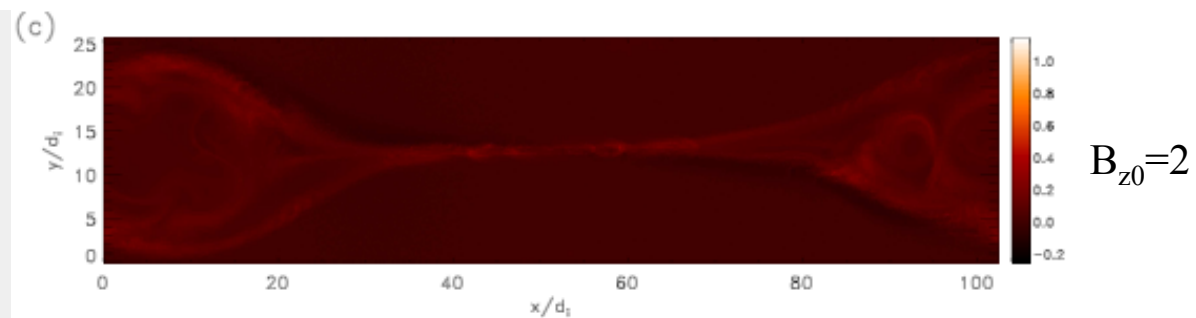
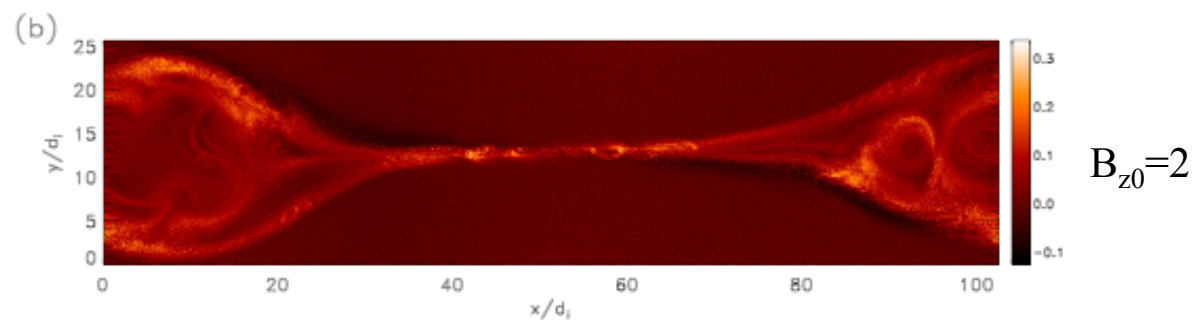
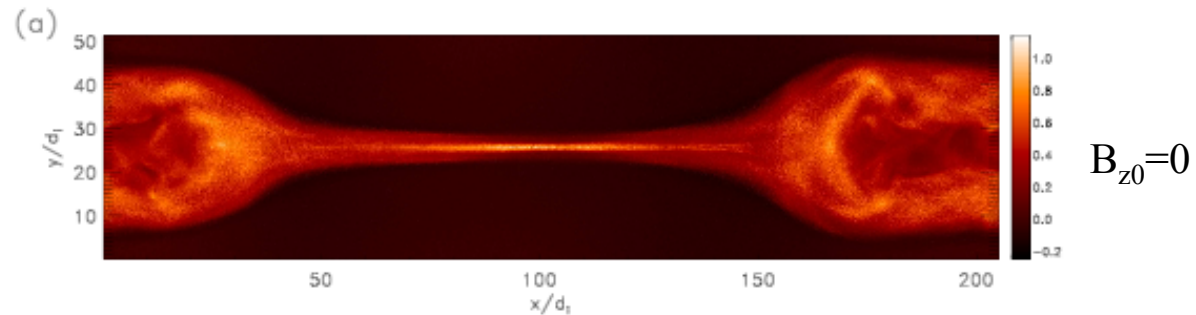
Seeding super-Alfvénic ions through pickup in reconnection exhausts

- Ion heating is dominated by large-scale reconnection exhausts rather than the localized region around the x-line
- Ions moving from upstream cross a narrow boundary layer into the Alfvénic reconnection exhaust
- The ion can then act like a classic “pick-up” particle, where it gains an effective thermal velocity equal to the Alfvénic outflow $T_i \sim m_i c_{Ax}^2$
 - during guide field reconnection there is a threshold for pickup behavior and heating is mostly in T_{\perp}

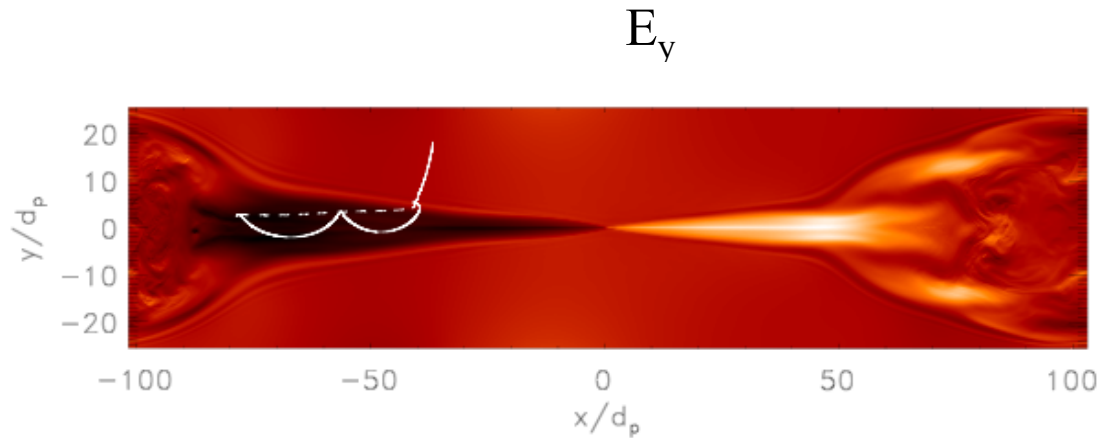
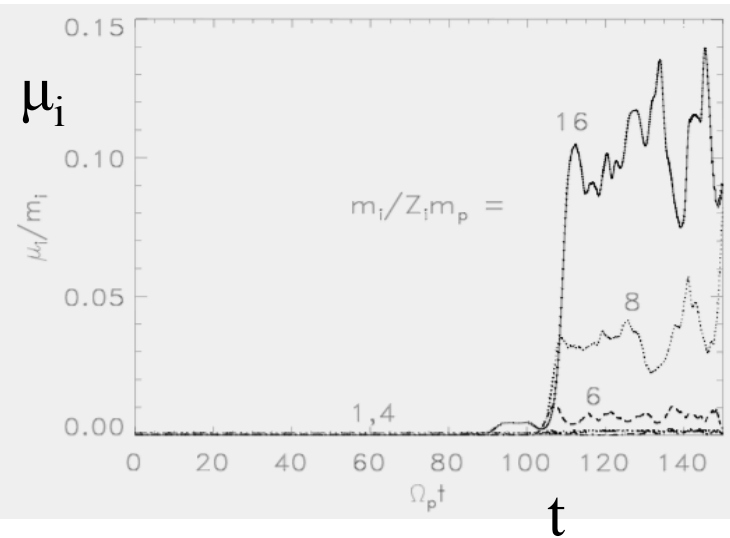


Ion temperature in reconnection outflows: anti-parallel versus guide field

- Comparison of PIC simulations with and without a guide field
- Temperature increments of protons
 - Little proton heating with strong guide field
 - Protons are adiabatic
 - Why?



Pickup threshold: guide field



$$B_{z0} = 5.0$$

- Protons and alpha particles remain adiabatic (μ is conserved)
- Only particles that behave like pickup particles gain significant energy \rightarrow **threshold for pickup behavior**

$$\frac{v_{iy}}{\Delta} \approx \frac{0.1c_{Apx}}{\rho_{sp}} > \Omega_i \Rightarrow \frac{m_i}{Z_i m_p} > \sqrt{\beta_{px}}$$

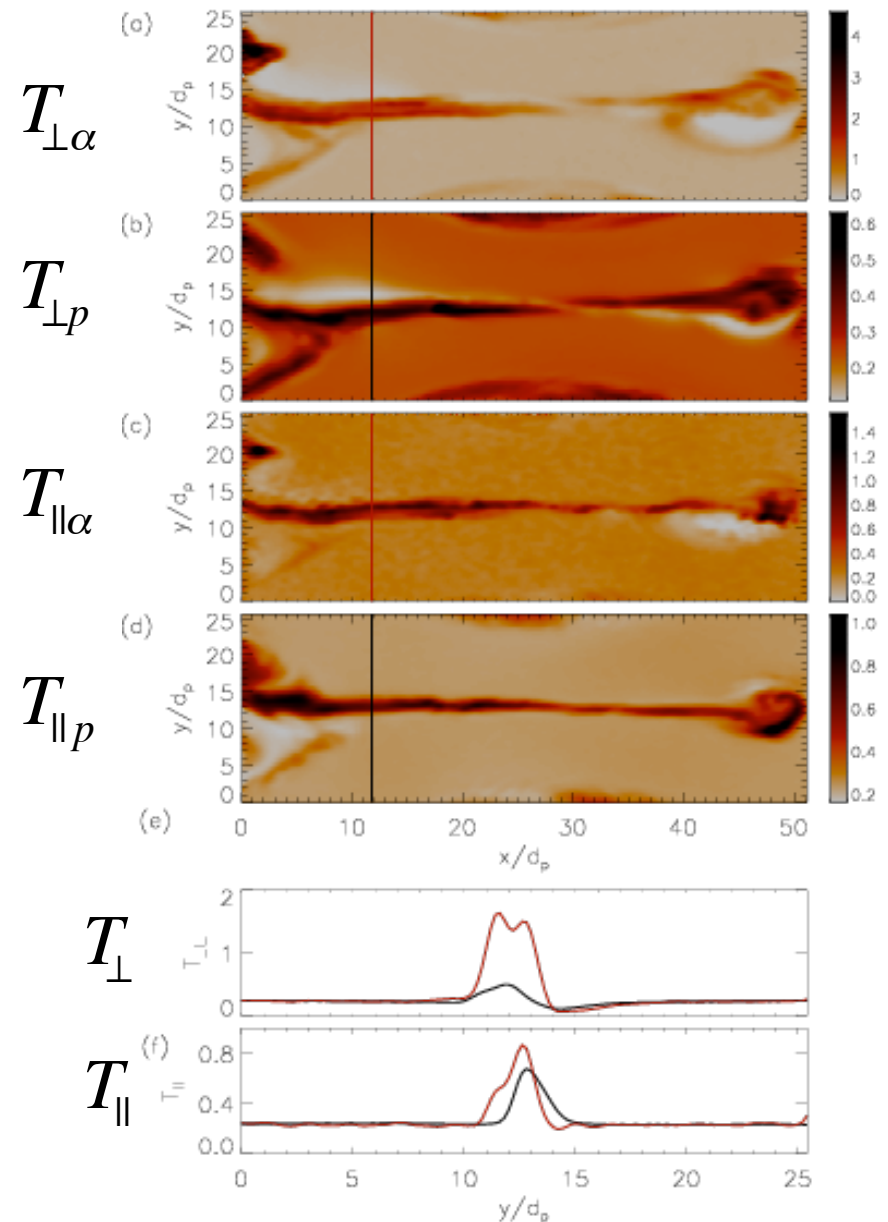
- Mostly perpendicular heating $\Delta T_{\perp} = \frac{1}{2} m_i c_{Ax}^2$ $\Delta T_{\parallel} = 0$

Reconnection with multiple ion species

- PIC simulations with a guide field 2.0 times the reconnecting field
 - Protons in the adiabatic regime
 - Include 1% fully stripped alpha particles
 - In the pickup regime

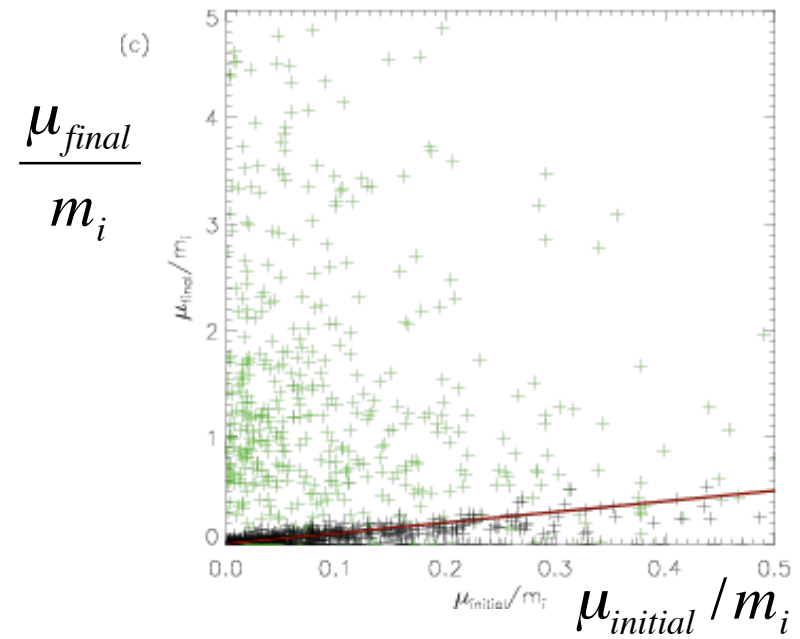
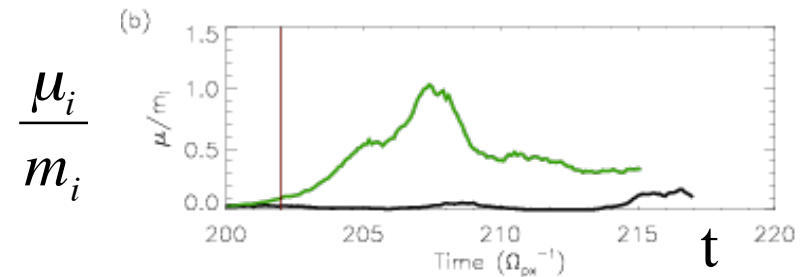
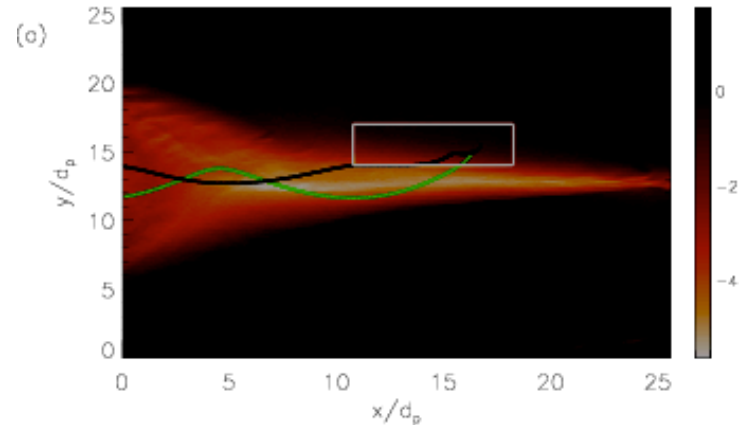
Alpha and proton heating

- Strong enhancement of $T_{\perp\alpha}$
 $T_{\perp\alpha} \gg T_{\parallel\alpha}$
 - Very different from anti-parallel reconnection
- Strong alpha heating compared to that of protons
 - Consistent with predictions



Pickup behavior of alphas

- Alpha trajectories are consistent with pickup behavior
- Strong increase in alpha magnetic moment μ_α
- No significant change in μ_p



Knizhnik et al '11

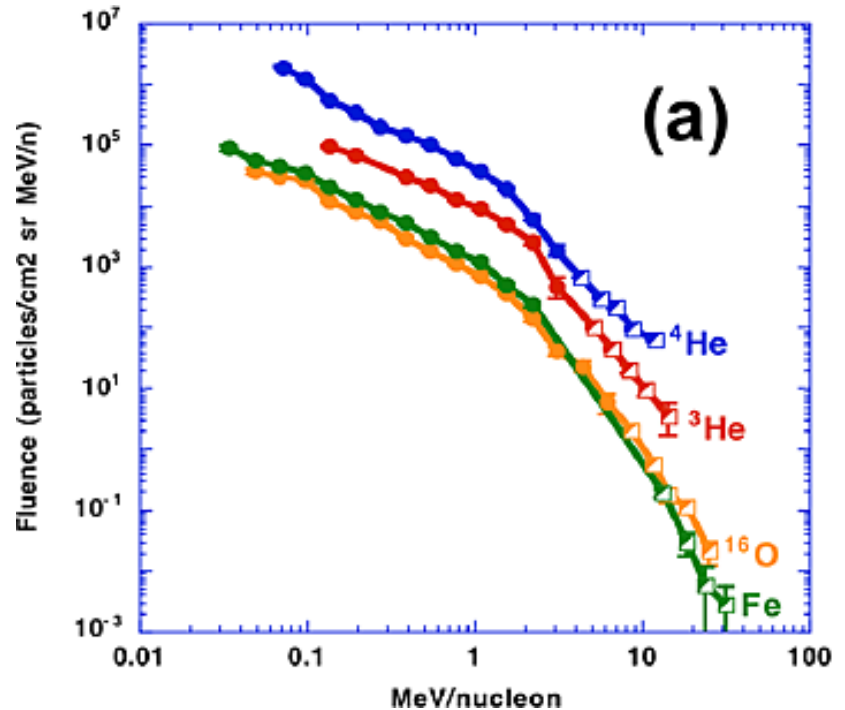
Reconnection exhaust heating

- For typical coronal parameters ($B \sim 50\text{G}$, $n \sim 10^9$), temperature increments $\sim 25\text{keV/nucleon}$
 - Typical of thermal component in flare heating
- Ion heating scenario
 - In a typically wide current sheet the reconnection magnetic field B_{0x} is very small $\Rightarrow \beta_{px} \sim 8\pi n T_p / B_{0x}^2 \gg 1$
 - Adiabatic behavior for all ions
 - As reconnection proceeds B_{0x} increases and β_{px} decreases and ions with progressively smaller $m_i / Z_i m_p$ behave like pickup particles and gain energy
$$\frac{m_i}{Z_i m_p} > \sqrt{\beta_{px}}$$
 - Mostly perpendicular heating
- Heavy trace ions gain energy first
- Consistent with coronal observations with $T_{\text{perp}} > T_{\parallel}$ and abundance enhancements in impulsive flares?

^3He abundance enhancement

- Abundance enhancements of ^3He range from $10^3 - 10^4$
 - Does not follow the typical M/Q trend of other ions
 - Spectra can differ significantly from other ions, especially for small flares
- Electron beam-driven ion cyclotron waves have been proposed as the heating mechanism for ^3He (Fisk '73, Temerin and Roth '96)
 - Volume of space with strong electron beams is limited because of constraints on the integrated current

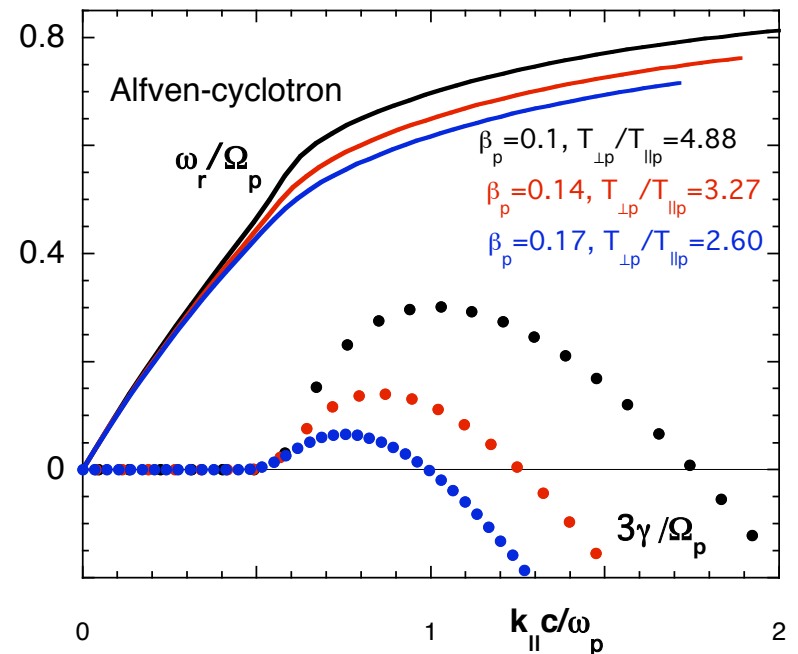
ACE ULEIS/SIS



Mason et al '02

Temperature anisotropy driven AIC waves

- The large ratio of T_{\perp}/T_{\parallel} in reconnection exhausts can drive Alfvén ion-cyclotron waves
 - Large volume of plasma with significant available free energy
- Are exploring this mechanism with PIC simulation with trace numbers of ^3He in a dominantly proton background with anisotropy

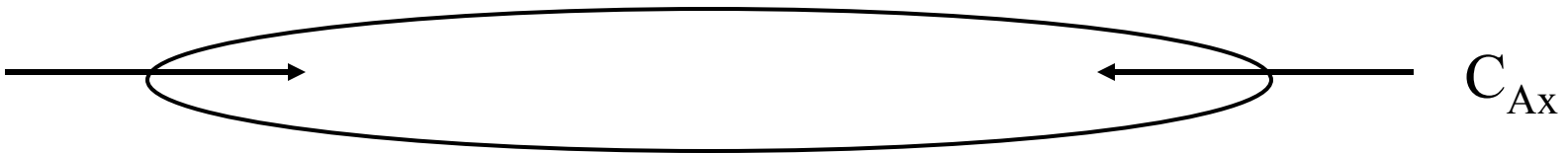


Gary et al '93

Particle acceleration in multi-island reconnection

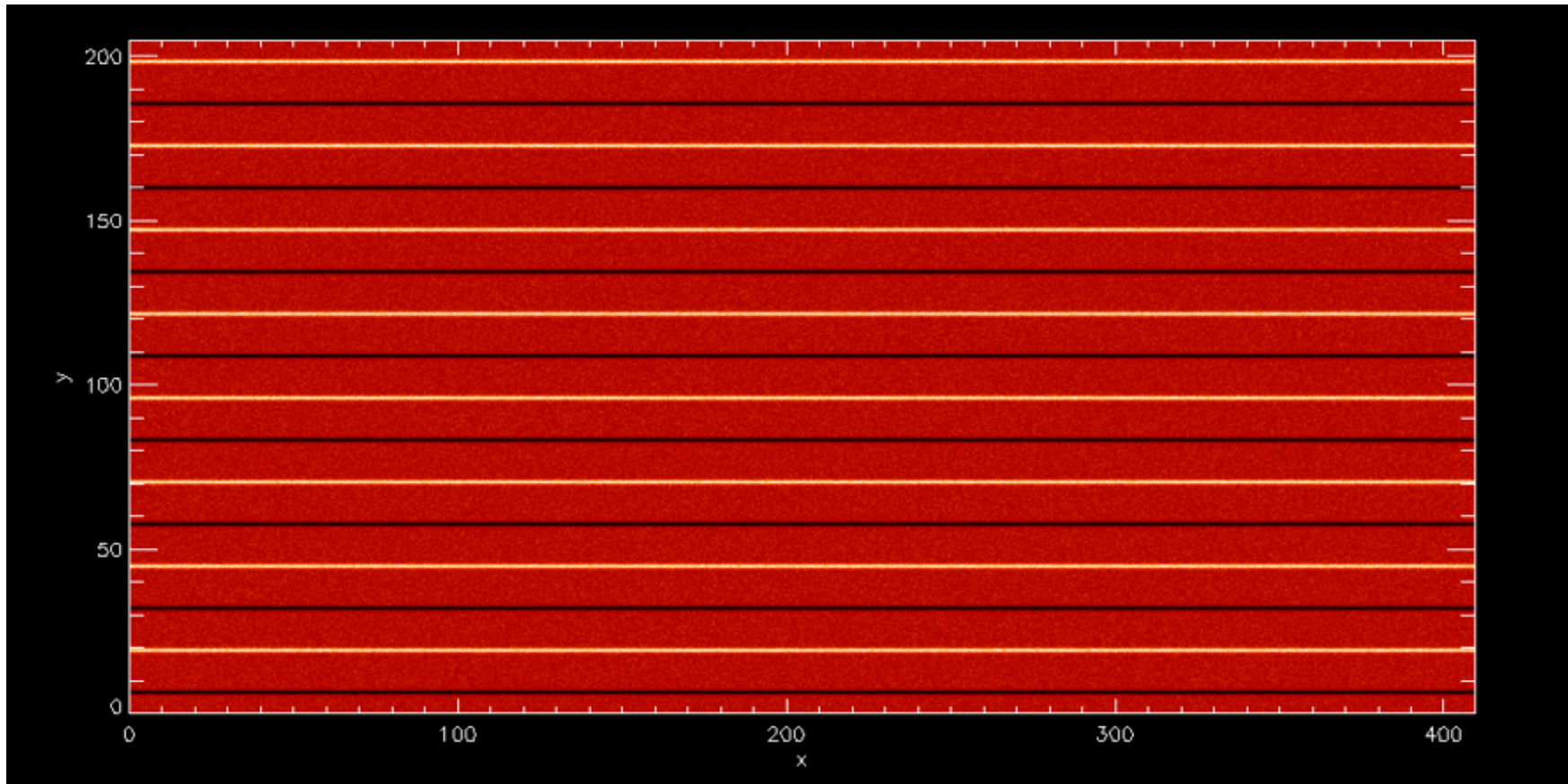
- Electrons and super-Alfvénic ions are accelerated to high energy in a multi-island environment
 - Fermi reflection in contracting magnetic islands (Kliem 94, Drake et al 2006, Drake et al 2010)
 - Dominantly parallel heating

$$\frac{d\varepsilon_{\parallel}}{dt} \sim 2\varepsilon_{\parallel} \frac{c_A}{L_x}$$



Particle acceleration during multi-island reconnection

- 2-D simulation in a multi-current layer system
 - Force-free with low initial β
 - Guide field equal to the reconnecting field



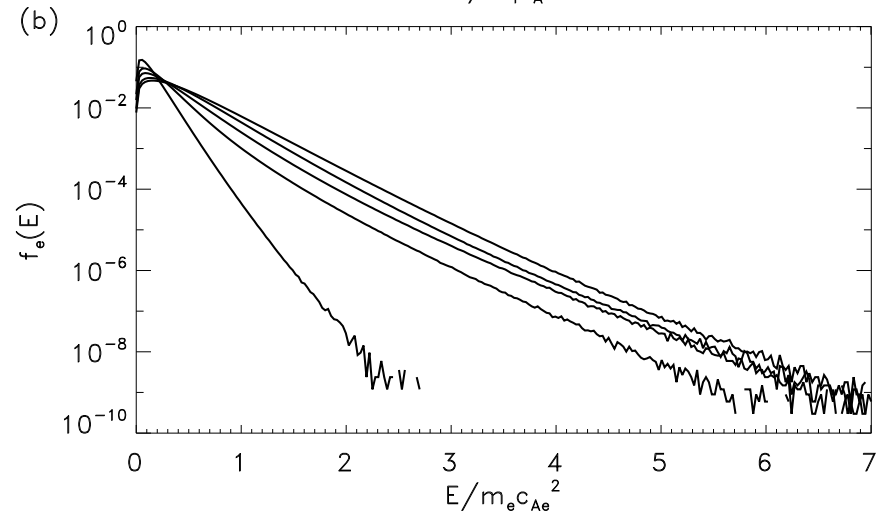
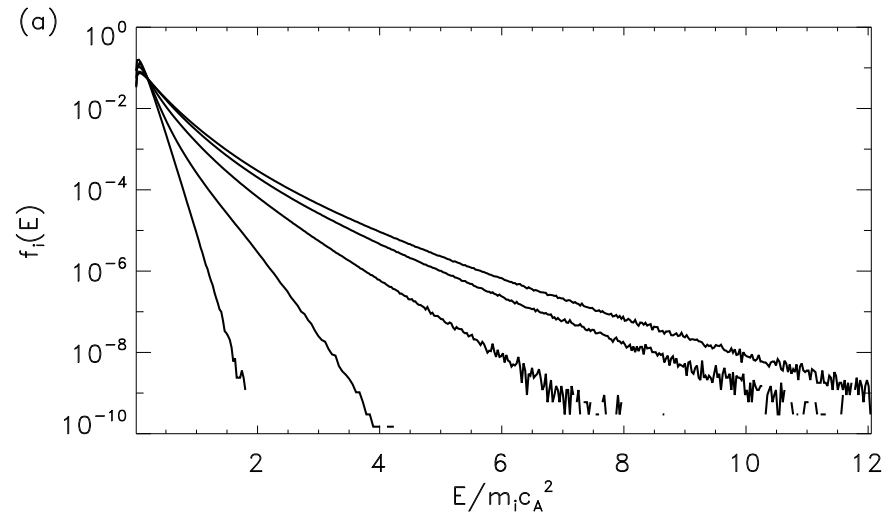
Electron and ion energy spectra

- Both ions and electrons gain energy
 - Electrons gain energy early and saturate. Why?
 - Ions gain energy later
- The rate of energy gain of particles increases with energy

⇒ consistent with first order Fermi

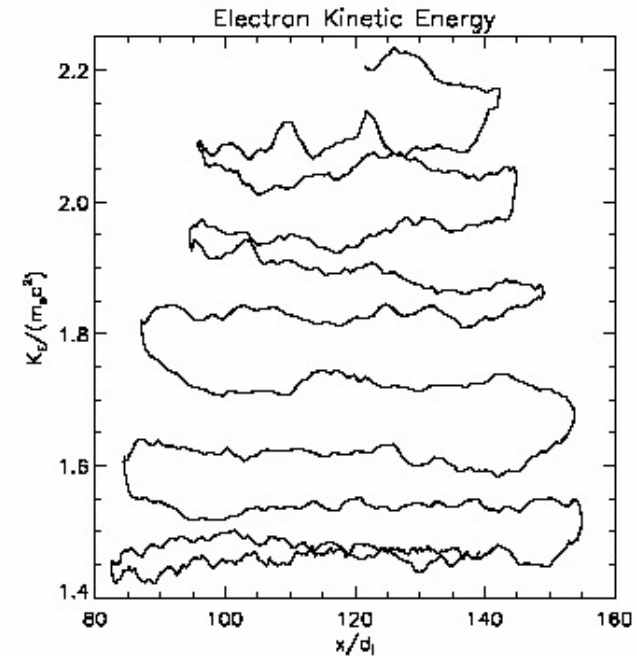
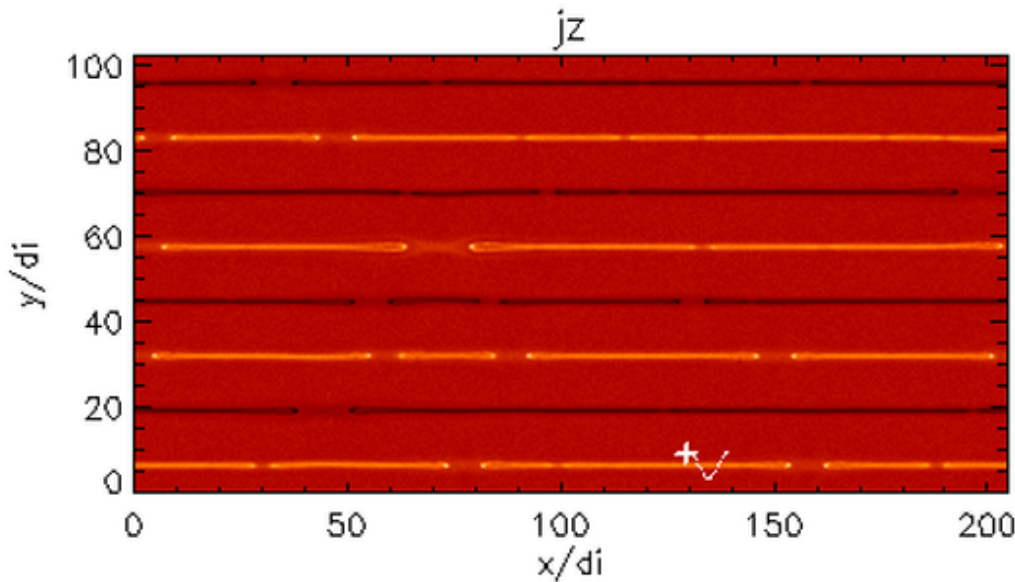
$$\frac{d\varepsilon}{dt} \propto \varepsilon$$

⇒ the simulation data is not a powerlaw since a log-linear plot



Fermi acceleration

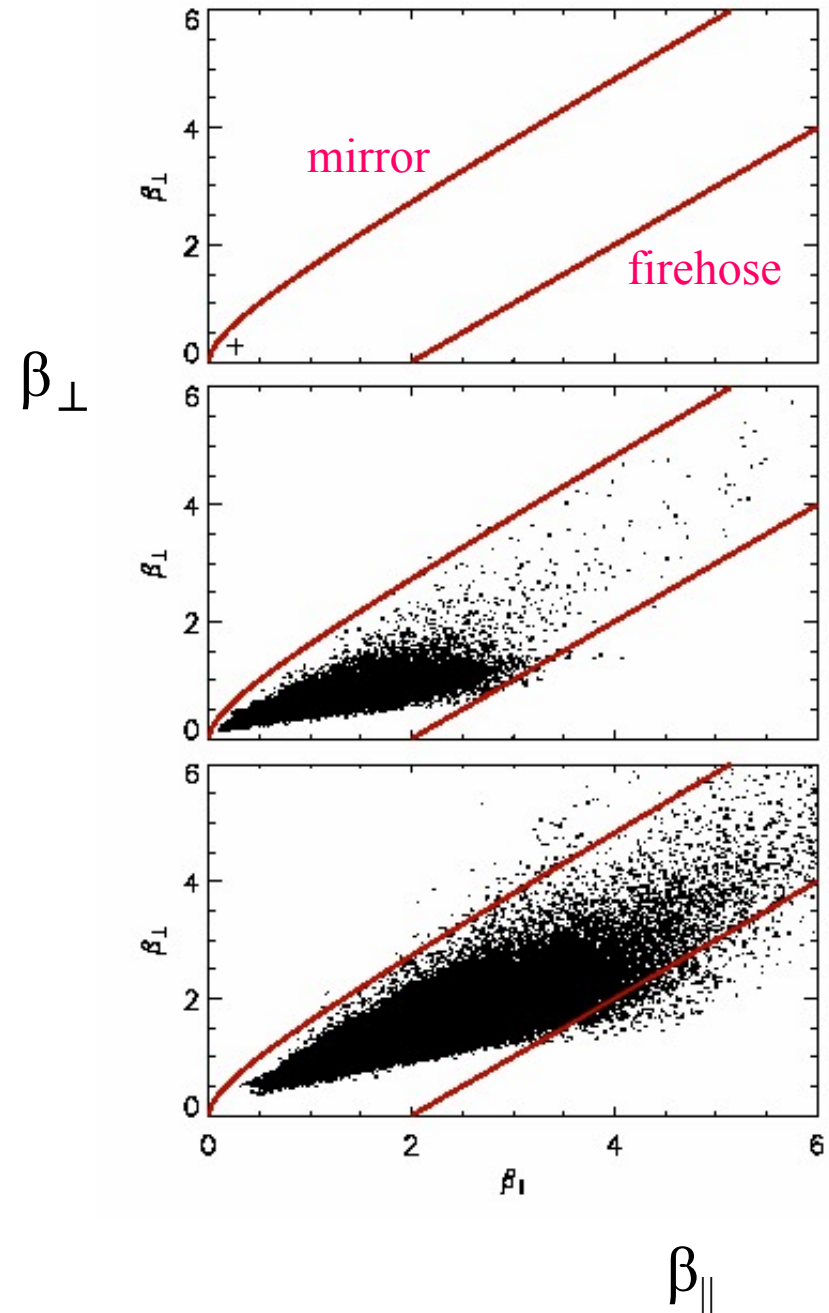
- How do the most energetic particles gain energy?
 - Reflection from the ends of contracting islands
 - Increase of parallel energy and pressure p_{\parallel}



$$\frac{d\varepsilon_{\parallel}}{dt} \sim 2\varepsilon_{\parallel} \frac{c_A}{L_x}$$

Mirror and firehose conditions: guide field

- Low initial β , guide field unity simulation
 - Each point corresponds to a grid point in the simulation
- At late time the islands bump against the firehose condition
- The energetic particles approach $\beta \sim 1$ even for low initial β
- Consistent with over-the-limb flare observations (Krucker et al 2010)



Energetic particle distributions

- Write down model equations for particle acceleration in contracting islands with feedback from the high pressure and convective loss
 - Powerlaw solutions for the omnidirectional energy flux

$$F(v) \sim \varepsilon^{-\gamma}$$

- Spectral index given by

$$\gamma = 1.5 + \frac{\beta_0}{2}$$

- Universal spectral index of 1.5 in low β systems

Conclusions

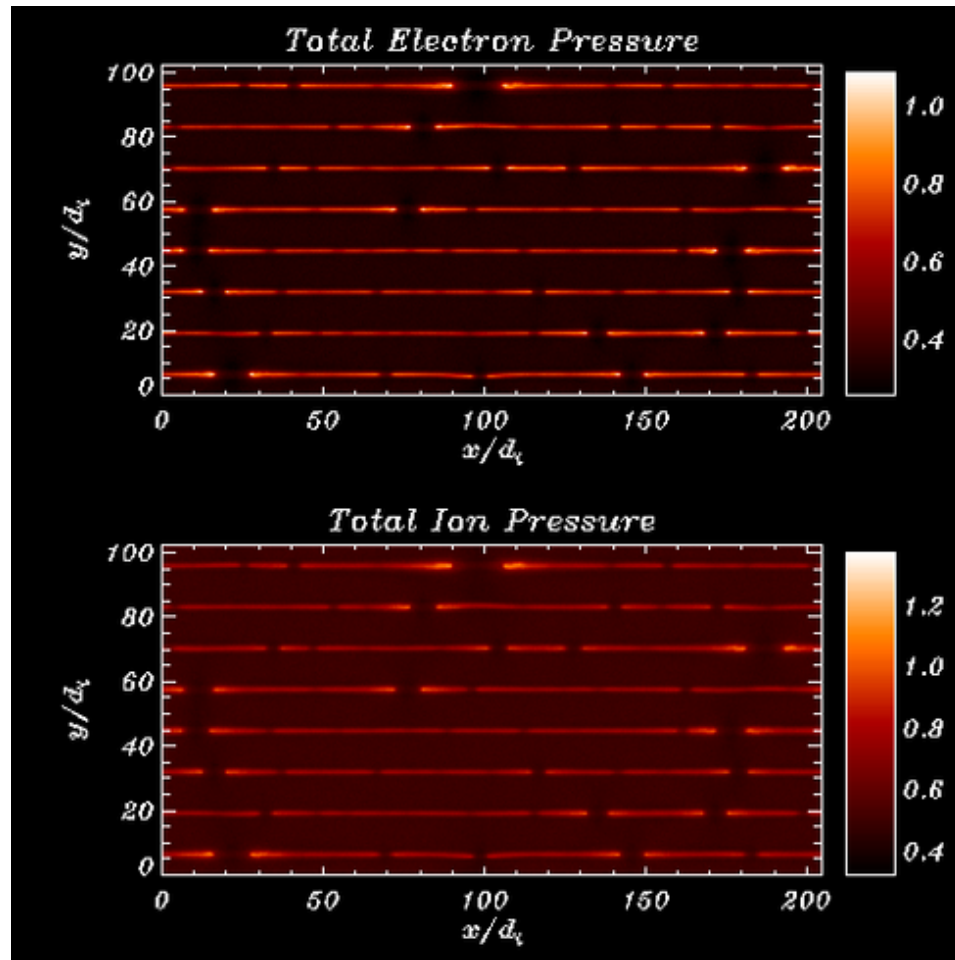
- Reconnection dominates reconnection in the solar wind close to the sun and in the corona where β is low.
- Ion interaction with the reconnection exhaust seeds them to super-Alfvenic velocities.
 - Ions act as pickup particles as they enter reconnection exhausts gain most energy
 - M/Q threshold for pickup behavior in guide field reconnection
 - Gain a thermal velocity given by the Alfven speed
 - Most of temperature increase is in T_{\perp}
- M/Q threshold for pickup behavior is a possible explanation of impulsive flare heavy ion abundance enhancements
- Alfven ion-cyclotron waves driven by $T_{\perp} > T_{\parallel}$ in reconnection exhausts preferentially heat ^3He and may be the source of abundance enhancements of ^3He

Conclusions (cont.)

- High energy particle production during magnetic reconnection requires the interaction with many magnetic islands
 - Not a single x-line
 - 1st order Fermi acceleration in contracting islands accelerates both ions and electrons
 - Island contraction is limited by the marginal firehose condition
 - Spectral indices of energetic particles take the form of powerlaws with spectral indices controlled by the firehose condition

Firehose instability during island contraction

- Fermi reflection within islands increases p_{\parallel} and leads to firehose



Schoeffler et al 2011

1-D Model equations

- Rate of energy gain: first order Fermi

$$\dot{v} = \frac{dv}{dt} = \frac{1}{\tau_h} \left(1 - \frac{4\pi p}{B^2} \right)^{1/2} v \quad \tau_h = \left\langle \frac{c_A}{L_w} \right\rangle^{-1}$$

Reduction of contraction rate due to firehose condition

- Model equation for the omnidirectional distribution function

$$F(v,t) = 4\pi v^2 f(v,t)$$

$$\frac{\partial F}{\partial t} + \frac{\partial}{\partial v} \dot{v} F = -\frac{1}{\tau_L} [F - F_0(v)] \quad \tau_L = \left\langle \frac{c_A}{L} \right\rangle^{-1}$$

- Above the source energy this is an equidimensional equation
 \Rightarrow powerlaw solutions