

Magnetic Reconnection with Optically Thin Radiative Cooling

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Conventional Reconnection Research:

- In conventional reconnection studies (space/solar/laboratory), the plasma consists charged particles (e -ns & ions) --- **no photons!**
- However, in many *astrophysical* situations, the energy density is so high that **radiation** should strongly affect the reconnection process:
 - Radiative cooling;
 - Radiation pressure;
 - Compton-drag resistivity.

Reconnection with (optically-thin)

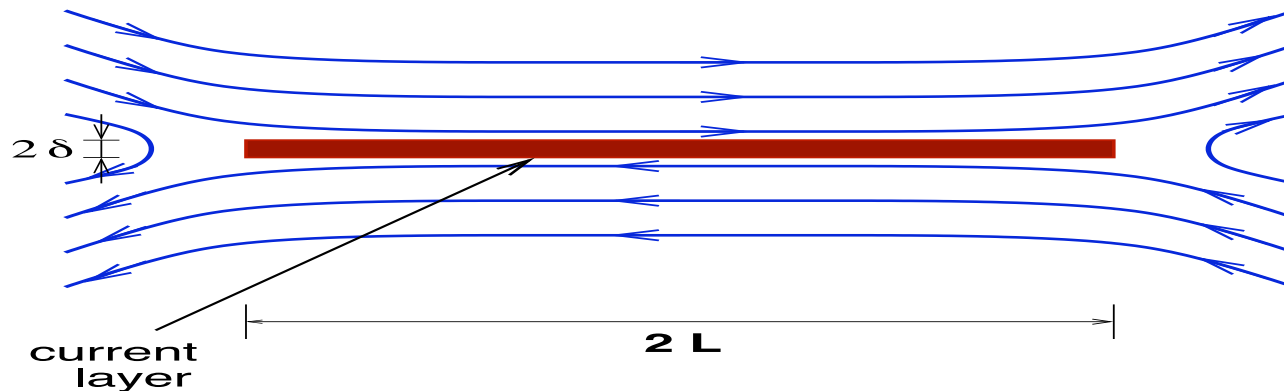
radiative cooling:

a Sweet-Parker-like analysis

(D. Uzdensky & J. McKinney, PoP 2011)

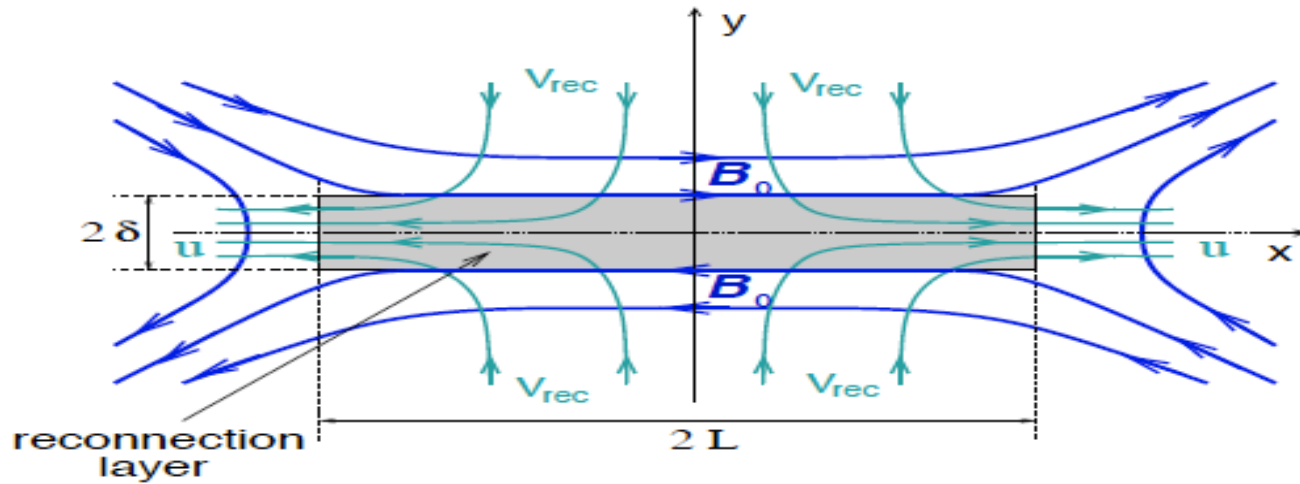
- Radiative cooling is the first radiation effect that comes into play as energy density is increased.
- Relatively easy to include.
- Sweet-Parker model is a natural first step in reconnection studies.

Two roles of radiative cooling



- Original SP model is incompressible -- OK even for low- β^{upstream} plasmas.
- In radiative reconnection, incompressibility gets replaced by heating/cooling balance, which then control layer temperature.
- Consequences of strong radiative cooling:
 - Strong plasma compression (if no guide field): $A = n/n_0 \gg 1$.
 - Lower $T_e \rightarrow$ higher resistivity \rightarrow faster reconnection.

Key parameters



- Upstream Alfvén velocity:

$$V_{A0} \equiv \frac{B_0}{\sqrt{4\pi n_0 m_p}} .$$

- Global Alfvén transit time:

$$\tau_{A,0} \equiv \frac{L}{V_{A,0}}$$

- Lundquist number:

$$S_0 \equiv \frac{LV_{A,0}}{\eta} \gg 1$$

Equations I:

- Mass Conservation: $n_0 v_{\text{rec}} L = n u \delta \Rightarrow v_{\text{rec}} L = A u \delta .$
- Faraday's law: $\partial_t B = 0 \Rightarrow \nabla \times E = 0 \Rightarrow E_z = \text{const}$
- Ohm's law: $c E_z^{\text{upstream}} = -v_{\text{rec}} B_0 = c E_z(0,0) = \eta' j_z(0,0)$
- Ampere's law: $j_z(0,0) \approx c \frac{B_0}{4\pi\delta}$
- Ohm + Ampere: $\eta = v_{\text{rec}} \delta$

Equations II: equations of motion

- Pressure balance across layer: $\frac{B_0^2}{8\pi} = P = 2nk_B T$

- Equation of motion along layer:

$$\rho v_x \partial_x v_x = -\partial_x p + j_z B_y / c.$$

- Reconnected magnetic field:
(from $cE_z = v_{\text{rec}} B_0 = u B_1$)

$$B_1 = B_0 \frac{v_{\text{rec}}}{u} \simeq B_0 \frac{\delta}{L} A$$

- Magnetic tension \gg pressure gradient:

$$j_z B_y / c \sim A B_0^2 / 4\pi L \gg \partial P / \partial x \sim B_0^2 / 8\pi L$$

- Outflow velocity: $u \sim \frac{B_0}{\sqrt{4\pi n_0 m_p}} = V_{A0}$

Solution in terms of A :

- Compression ratio:

$$A \equiv \frac{n}{n_0}$$

(n_0 – upstream density; n -- density in the layer)

- Reconnection velocity: $v_{\text{rec}} = V_{A0} S_0^{-1/2} A^{1/2}$

- Layer thickness: $\delta \sim \frac{\eta}{v_{\text{rec}}} \sim L S_0^{-1/2} A^{-1/2} = \delta_{\text{SP}} A^{-1/2}$

Equations III: Heating/cooling balance determines compression ratio A

- Ohmic heating rate: $Q_{\text{ohm}} = \eta' j^2 \sim A \frac{B_0^2}{4\pi\tau_{A,0}} \equiv A Q_0$
where $Q_0 \equiv B_0^2 / (4\pi\tau_{A,0})$

- Advective cooling rate: $Q_{\text{adv}} \sim \frac{nkT}{L/u} \sim \frac{B_0^2}{4\pi\tau_{A,0}} \simeq Q_0$

- Radiative heating/cooling balance:

$$Q_{\text{rad}}(n, T) = Q_{\text{ohm}} \sim A Q_0$$

- Equation for A:

where

$$\begin{cases} n(A) = n_0 A \\ k_B T(A) = k_B T_{\text{eq}} A^{-1} \equiv A^{-1} \frac{B_0^2}{16\pi n_0} \end{cases}$$

$$Q_{\text{rad}}[n(A), T(A)] = \frac{B_0^2}{8\pi\tau_{A,0}} A = A Q_0$$

Conditions of validity

- Strong cooling/strong compression: $A \gg 1$

- Evolutionary condition:

plasma should be able to cool down to this state

$$Q_{\text{rad}}[n_0, T_{\text{eq}}(n_0)] > Q_{\text{Ohm}}[n_0, T_{\text{eq}}(n_0)] \simeq Q_0$$

- Combining with $Q_{\text{rad}}(A) = Q_{\text{ohm}}(A) \simeq A Q_0$,

we get a condition on the cooling function:

$$Q_{\text{rad}}[n(A), T(A)] < A Q_{\text{rad}}[n_0, T_{\text{eq}}(n_0)]$$

Which radiative processes allowed?

- Evolutionary condition:

$$Q_{\text{rad}}[n(A), T(A)] < A Q_{\text{rad}}[n_0, T_{\text{eq}}(n_0)]$$

- Consider general form:

$$Q_{\text{rad}} \sim n^\alpha T^\beta$$

- Then, strong radiation regime requires: $\alpha < 1 + \beta$

- Important special cases:

- Bremsstrahlung ($\alpha=2, \beta=1/2$)
- Cyclotron ($\alpha=1, \beta=1$)
- Inverse Compton ($\alpha=1, \beta=1$)

Inverse Compton (IC) Cooling

- Important for accreting black hole coronae/jets, etc.

- IC cooling rate per electron: $Q_{\text{rad}}/n = \frac{4}{3} \sigma_T c U \beta^2 \gamma^2$

where $U = U_{\text{rad,ext}}$

- Non-relativistic limit: $Q_{\text{rad}} \sim \sigma_T c U \frac{n T_e}{m_e c^2} \sim \sigma_T c U \frac{P}{m_e c^2}$

- Compression ratio: $A = \frac{Q_{\text{rad}}}{Q_0} \sim \sigma_T c \frac{U \tau_{A,0}}{m_e c^2}$

Cyclotron/Synchrotron Cooling

- Cyclo/synchrotron cooling rate per electron:

$$Q_{\text{rad}}/n = \frac{4}{3} \sigma_T c U \beta^2 \gamma^2$$

where $U = U_{\text{magn}} = B^2/8\pi$

- Nonrelativistic limit:

$$Q_{\text{rad}} \sim \sigma_T c U \frac{n T_e}{m_e c^2} \sim \sigma_T c U \frac{P}{m_e c^2}$$

- Compression ratio:

$$A = \frac{Q_{\text{rad}}}{Q_0} \sim \sigma_T c \frac{U \tau_{A,0}}{m_e c^2} \sim \frac{m_i}{m_e} \frac{V_{A,0}}{c} \tau_T$$

where

$$\tau_T = \sigma_T n_0 L$$

Work in Progress/Open Questions:

- Generalization to finite guide field
- Resistive-MHD validity (collisionality): $\delta > d_i, \rho_i$
- Electron-ion thermal equilibration: $T_e = T_i$?
- Generalization to finite optical depth
- Radiation pressure, Compton resistivity
- Astrophysical applications
- HED Laser-Plasma laboratory experiments

(Nilson et al. 2006; C.K. Li et al. PRL 2007)

SUMMARY

- Radiative effects:
 - radiative cooling;
 - radiation pressure;
 - radiative (Compton) resistivity

are important in high-energy-density astrophysical environments.

- We started a program of investigating the effects of radiation on reconnection.
- First step: Sweet-Parker theory for resistive reconnection w. radiative cooling.
- Intense radiative cooling:
 - lowers layer temperature, which makes it more collisional and increases resistivity and hence reconnection rate;
 - unless there is strong B_{guide} , cooling + pressure balance \rightarrow strong plasma compression, which also makes reconnection faster.
- Several astrophysically-important radiative mechanisms (bremsstrahlung, cyclo-synchrotron, and inverse-Compton) explored.