





Crystallographer as seen by others

Crystallographer as seen by myself



















What d	o we w QC s	/ant to know about a structure?
Crystallographer:	orderin structu	g principles of idealized and real real
		As far as possible model-free structure analysis in order to find out whether or not QCs are quasiperiodic and what kind of disorder is present. One of the goals is to provide realistic models as input for quantum- mechanical calculations.
Structure based	understand	ling of stability and properties of matter

	What do	we want to know about a QC structure?
	We want to know ev describes reality rea	erything we need to build a model that sonably well.
٢	Mathematician:	idealized geometrical-structure building principles - mathematical objects
	Physicist:	idealized factors governing structure formation and physical properties
ŀ	Materials Scientist:	real structure/microstructure-property relationships
L	Crystallographer:	ordering principles of idealized and real structures
	Structure based unde	rstanding of stability and properties of quasicrystals

Structural peculiarities of quasicrystals

Long-range order

- > Are QC strictly quasiperiodic or on average only or....
- Is nD description really applicable? If yes, are atomic surfaces dense or fractal or ...
- How does long-range order vary with temperature, pressure ...

Disorder

- Does disorder in quasicrystals differ from that in crystals?
- What is the structural meaning of phasons and phason fluctuations?

/here are the atoms (each one) in a quasicrysta



Structural peculiarities of quasicrystals

Nowadays, ab-initio quantum-mechanical modeling is possible for approximants only (< 1000 atoms). Consequently, the property that is typical for quasicrystals, *i.e.* the quasiperiodic long-range order, cannot be studied in this way at present and near future.

Where are the atoms - long-range order



Size of full data sets









Structural peculiarities of quasicrystals

Short-range order (the structure of clusters)

The method of choice for obtaining a starting model are electron microscopic methods, in particular for decagonal phases.

The novel Ultra-High-Resolution Transmission Electron Microscopes (UHRTEM) with full correction of spherical aberration will allow to get much better images in future. For instance, for a 200 keV microscope (Zeiss), a resolution of 0.8 Å is possible.

However, compared to diffraction methods, the accuracy is very low (laterally: \approx 0.5 Å, vertically: no spatial resolution at all).

Determination of local order



Structural peculiarities of quasicrystals

Disorder

A principal problem is that in most cases non-equilibrium structures are studied. Usually, a QC is prepared from the melt, HT annealed, quenched.

During cooling to RT partial relaxation takes place due to a different efficient atomic volume at HT and RT (atomic vibrations). Chemical disorder, thermal vacancies and random phason fluctuations are not equilibrated.

Structural disorder of the displacive and/or the substitutional type may be present anyway, also at equilibrium conditions.

Structural disorder





















































STM only shows a kind of electronic charge distribution as a function of the tunneling current. The image obtained reflects the

electronic surface structure rather than the atomic structure

















Exar	nple NAD⁺ synthetase		
	P2 ₁ , a=52.28, b=84.97, c=59.64 Å, β=110.5°		
	No. of measured reflections	790 765	
	No. of unique reflections	231 200	
	with I > 4 σ (I)	170 606	
	No. of reflections in refinement	205 215	
	No. of parameters	47 101	
	No. of restraints	57 173	
	No. of non-H atoms (excl. waters)	4 518	
	R-factor	0.116	
Symers	ky et al, Acta Cryst. D58 (2002) 1138	59	



Average-structure solution by MEM
Solving decagonal structures requires the fundamental maximum entropy equations
(Bricogne 1984)

$$q_i = \frac{p}{Z(\lambda_1,...,\lambda_N)} \exp\left(\sum_{n=1}^{N} \lambda_n \frac{\partial C_n}{\partial q_i}\right)$$
(5)

$$Z\left(\lambda_1,...,\lambda_N\right) = \sum_{i=1}^{N} p \exp\left(\sum_{n=1}^{N} \lambda_n \frac{\partial C_n}{\partial q_i}\right)$$
(6)
to be solved in five dimensions for all N_P grid points q_i restricted to all N_C constraint
equations C_n by Lagrange multipliers $\dot{\chi}_n$. Two different constraint equations are
necessary to take all structure factors derived from the symmetry minimum solution and
all the observed structure amplitudes simultaneously into account. Assuming the noise to
be gaussian the known structure factors F_{ebb} (H) can be constrained by

$$C_{1} = \sum_{H} \frac{1}{\sigma^{2}} \left| F_{obs} \left(H \right) - F_{clc} \left(H \right) \right|^{2} = \chi^{2}.$$
(7)

Haibach et al., Acta Crystallogr. A (1996) 277

Average structure solution by MEM
with their corresponding standard deviations
$$\sigma$$
. The second constraint equation (compare
Sakata & Sato 1990) only depends on the structure amplitudes and restrains all unknown
phases to

$$C_{2} = \sum_{\mathbf{x}} \frac{1}{\sigma^{2}} \left| \left| F_{abc}(\mathbf{K}) \right| - \left| F_{cb}(\mathbf{K}) \right| \right|^{2} = \chi^{2}.$$
(8)
Substituting (7) and (8) into equation (5) results in a five-dimensional algorithm

$$q_{i} = \frac{\mu}{Z\left(\lambda_{i},\lambda_{2}\right)} \exp\left[-2\lambda_{i}\sum_{\mathbf{H}} \frac{1}{\sigma^{2}} \left| F_{abc}(\mathbf{H}) - F_{cb}(\mathbf{H}) \right| \cos\left(2\pi H\mathbf{x}_{i} - \varphi_{\Lambda}\right) - 2\lambda_{2}\sum_{\mathbf{k}} \frac{1}{\sigma^{2}} \left| \left| F_{abc}(\mathbf{K}) \right| - \left| F_{cb}(\mathbf{K}) \right| \left| \cos\left(2\pi K\mathbf{x}_{i} - \varphi_{\Lambda}\right) \right|,$$
(9)

with $\phi_{\Delta} = \arctan \left\{ \ln \left[F_{obs}(H) - F_{obs}(H) \right] \right\} / Re \left[F_{obs}(H) - F_{obs}(H) \right] \right\}$, the phase of the structure factor difference. This equation can be maximised solving λ by Newton's method (Bricogne 1984) or by exponential modelling (Collins & Mahar 1983).

Haibach et al., Acta Crystallogr. A (1996) 27









Structure refinement - average structure

SHELXL always refines against F². Refinement against ALL F²-values is demonstrably superior to refinement against F-values greater than some threshold [say 46(F)]. More experimental information is incorporated (suitably weighted) and the chance of getting stuck in a local minimum is reduced. In pseudo-symmetry cases it is very often the weak reflections that can discriminate between alternative potential solutions. It is difficult to refine against ALL E-values because of the difficult yold estimating $\alpha(F)$ from $\alpha(F^2)$ when F^2 is zero or (as a result of experimental error) negative. The diffraction experiment measures intensities and their standard deviations, which after the various corrections give F_0^2 and $\sigma(F_0^2)$.

The use of a threshold for ignoring weak reflections may introduce bias which primarily affects the atomic displacement parameters; it is only justified to speed up the early stages of refinement. In the final refinement ALL DATA should be used except for reflections known to suffer from systematic error. Anyone planning to ignore this advice should read Hirshfeld & Rabinovich (1973) and Arnberg, Hovmöller & Westman (1979) first. Refinement against F² also facilitates the treatment of twinned and powder data, and the determination of absolute structure.

owchart

R-factors and Goodness of fit



where n is the number of reflections and p is the total number of parameters refined. Weighting scheme $~w=1~/~[~\sigma^2(F_o^{-2})~].$

One cosmetic disadvantage of refinement against F^2 is that R-indices based on F^2 are larger than (more than double) those based on F.

The R-index for MEM calculations is not a reliability factor. Its value is meaningless.





Diffraction methods - weak reflections

Weak reflections in periodic structure analysis

Due to atomic scattering factor and temperature factor large q^{II}-reflections are weak. They are important for high-resolution electron density maps but not for atomic positions.

Weak reflections in quasiperiodic structure analysis

Additional to weak reflections of the type described above, all reflection intensities are rapidly falling off with q^{\perp} independent from $q^{||}.$

Problem

All diffraction experiments on high-quality single crystals suffer from multiple diffraction effects, which may strongly increase intensities of weak reflections. This is a special problem for quasicrystals with their dense reciprocal space.



X-ray data collectio	on - diffuse intensities
3D diffuse diffraction data ca area detectors and synchrotic help distinguishing between	an only be quantitavely collected employing ron radiation. Variation with temperature can static and dynamic disorder.
Separation of diffuse scattering	ing of static and dynamic origin by quasielastic g.
Problem: Separation of Bragg reflection	n intensities from diffuse scattering
ocparation of Dragg Tenectio	in interfactors from diffuse southering.









Modelii	ng based o	on first-p	orincipl	es		
					-	
Based on methods v properties	realistic structur vill give realistic	al paramete results for th	rs modeling ne relaxed :	g by first-prir structure and	nciples d physical	
Problem:						
<i>Problem:</i> Till now, c atoms per	nly approximant unit cell < 1000	s can be cal	Iculated wit	h a maximu	m number	of
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Where will we be at ICQ10?

We will certainly have toy models, which much better describe experimental observations than actual models do, especially for decagonal Al-Co-Ni.

We will better understand what clusters are and what properties they have. Consequently, we will better understand the local (short-range) order in quasicrystals.

We will still not have solved the long-range order structure of quasicrystals with an accuracy comparable to that of regular structure analysis.

Let's start working!