# A Stable and Accurate Method for Tetrahedral Elastic-Plastic Computations

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Solids on Tets

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# Motivation: Tetrahedral Meshes

- Solid mechanics on hex meshes: mixed staggered Q1/Q0 formulation
  - continuous linear kinematic variables
  - discontinuous piece-wise constant stresses
  - requires various hourglass controls (e.g. Belytschko-Flanagan)
  - mesh generation very time consuming
- Tet meshes for solids:
  - use of automated fast meshing
  - ease of use for mesh adaptivity
  - ease of coupling with other physics (thermal, electromagnetic)

#### **Overview of Recent Research**

- Swansea: Bonet, Burton, Marriot, Hassan, (P1/P1-projection)
- Sandia: Dohrman, Key, Heinstein, Bochev, (P1/P1-projection)
- TU Munich-Sandia: Gee, Wall, Dohrman, (P1/P1+P1/P0-+proj.)
- LLNL: Puso, Solberg, (P1/P1+P1/P0-+proj.)
- RPI: Maniatty, Klaas, Liu, Shephard, Ramesh, (P1/P1-stabilized)
- Chorin's projection: Onate, Rojek, Taylor, Pastor (P1/P1)
- UPC Barcelona: Chiumenti, Cervera, Valverde, Codina (P1/P1-stabilized)
- UIUC: Nakshatrala, Masud, Hjelmstad, (P1/P1+bubble)
- Swansea II: Bonet, Gil, (P1/P1-stabilized)
- Berkeley/Pavia: Taylor, Auricchio, Lovadina, Reali, (Mixed enhanced)
- UCSD/University of Padua: Krysl, Micheloni, Boccardo (Mixed enhanced)
- Caltech: Thoutireddy, Ortiz, Molinari, Repetto, Belytschko (Composite Tets)

# Finite Elements: Fluids and Solids

- Our approach to solids is an extension of the VMS-stabilized hydro approach <sup>1</sup>
- All variables are nodal except deviatoric stress and internal material state variables, which are based at quadrature points
- P1-based tets enable use of one-point quadrature (as in uniform gradient hexes)



<sup>1</sup>G. Scovazzi, J. Comput. Phys., Vol 231 (24), 2012, pp. 8029<del>\_</del>8069<u></u> → ( = → ) = → ) へ

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# Governing Equations (Mixed Form)

 Solve for {d, v, σ̄, p} satisfying mass/momentum conservation, Cauchy stress decomposition, and velocity definition:

$$\rho J = \rho_0, \quad \rho \dot{v} = \nabla \cdot \sigma + \rho \cdot b, \quad \sigma = p I + \bar{\sigma}, \quad d = v.$$
(1)

- Assume σ
   is a function of the kinematics (strains, strain rates), state variables, the history of σ
   , etc.
- In the linear case we consider the **mixed** system for displacement
   (u) and pressure (p):

$$\rho \ddot{u} - \nabla \cdot \overline{\varepsilon}(u) - \nabla p = f$$
$$p - \kappa \nabla \cdot u = 0$$

where  $\overline{\varepsilon}(u)$  is the deviatoric strain tensor.

# Linear Elasticity: Static Case

- Stabilization for linear elasticity is very similar to Stokes flow
- Incompressible case:
  - P1/P0 locking (as in P1 displacement formulation)
  - P1/P1 checkerboard instability for pressure
- Solution for P1/P1: Hughes/Franca/Balestra stabilization (1986): enrich the velocity/displacement (*u*) with a residual-based term

$$u = u_h + u', \quad u' = -\tau \frac{h^2}{2\mu} (-\nabla p_h - \nabla \cdot \epsilon(u_h) - f)$$

- Stabilization derives from the additional pressure Laplacian
- This is now called Variational Multiscale (VMS) stabilization

# Linear Elasticity: Dynamic Case

- The Hughes/Franca/Balestra stabilization extends naturally to time-dependent Stokes/Navier-Stokes flows
- We could not find an appropriate  $\tau$  that worked for linear dynamics
- The issue appears to be the different character of the PDEs:
  - elliptic: Stokes (velocity/pressure)
  - elliptic: elasticity (displacement/pressure)
  - parabolic: time-dependent Stokes
  - hyperbolic: time-dependent elasticity
- Our solution is to formulate the pressure equation in rate form: which pairs naturally with the momentum equation:

$$\kappa^{-1}\dot{p} - \nabla \cdot v = 0$$
  

$$\rho \dot{v} - \nabla p = \nabla \cdot \overline{\varepsilon}(u) + f$$

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# Linear Elasticity: Dynamic VMS

- The stabilization then is analogous to what is used for the linear acoustic wave equation
- We add in subgrid scales  $\{v', p'\}$  defined using residuals

$$v = v_h + v', \quad v' = -\tau \rho^{-1} (\rho \dot{v}_h - \nabla p_h - \nabla \cdot \bar{\epsilon}(u_h) - \rho \cdot b)$$
  
$$p = p_h + p', \quad p' = -\tau (\dot{p}_h - \kappa \nabla \cdot v_h)$$

- The resulting pressure Laplacian and velocity div-div terms provide stabilization
- The use of residuals provides consistency and thus accuracy

# Linear Elasticity: Verification

- We verified the linear elastic case under various options:
  - static/dynamic,
  - quad/tri/hex/tet,
  - compressible/nearly incompressible,
  - structured/unstructured grids
- Plots of manufactured solutions for pressure:



Compressible (left) and nearly incompressible (right)

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## Linear Elasticity: Verification

- Verification test: analytic pressure/velocity/displacement with valid solution in the incompressible limit (here ν = 0.4995)
- Expected convergence rates are order 2 except for pressure which is order 1.5.



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## Nonlinear Dynamics: Hyper-elasticity

• We concentrate on mixed formulations using a pressure:

$$\sigma = pI + \bar{\sigma}$$

• Pressure is assumed a function of the volumetric part of the deformation gradient:

$$p \equiv \kappa \, U'(J) \tag{2}$$

where J = det(F), *F* is the deformation gradient, *U* is an energy function and  $\kappa$  is the bulk modulus (e.g.  $U(J) = \frac{1}{2}(J-1)^2$ )

• Deviatoric stress is defined in terms of J and  $b = F F^T$  using another energy function, for example using a neo-Hookean law

$$\bar{\sigma} = \mu J^{-5/3} \, \bar{b}$$

# Nonlinear Dynamics: J<sub>2</sub> Plasticity

- For plasticity, the pressure often remains a function of J
- The deviatoric stress is computed through an associative flow rule, with inclusion of constraints and plastic strain
- We have implemented a simple plasticity model <sup>2</sup> using linear hardening and a product factorization of the total deformation:

$$F = F^e F^p$$

- Extensions to other models (e.g. hypo-elasticity) should be possible provided that
  - we have separate models for  $\bar{\sigma}$  and p, and
  - the pressure remains a function only of J

<sup>&</sup>lt;sup>2</sup>JC Simo, CMAME (1992), pp. 61-102

## Pressure Evolution Equation and VMS

• The nonlinear pressure equation in an evolution form:

$$\dot{p} = \frac{\partial}{\partial t} \{ \kappa U'(J) \} = \kappa U''(J) \dot{J}$$
$$= \kappa U''(J) \{ J \nabla \cdot v \}$$
$$= \tilde{\kappa}(J) \nabla \cdot v.$$

• We have defined an effective bulk modulus that varies as the material undergoes volume change.

$$\tilde{\kappa}(J) \equiv \kappa \, U''(J) \, J$$

• VMS stabilization proceeds as in the linear case using a v' term

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# Nonlinear Examples Using Tet Meshes

#### Dynamic

- Taylor bar impact (elastic-plastic)
- Bending beam (hyper-elastic)
- Quasistatic (run in dynamic mode)
  - Billet in compression (elastic-plastic)
  - Cylindrical bar in uniaxial tension (elastic-plastic)
- Impact test in complex geometry

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# **Taylor Bar**

- Length/Radius: 3.24/0.32cm, density: 8930
- Elastic-plastic material (E=117.0e9,  $\nu$ =0.35,  $\sigma_{\gamma}$ =0.4e9, H=0.1e9)
- Zero normal velocity at wall, initial velocity 227m/s



# Taylor Bar: Pressure Convergence



Note: unstructured grids fill space more evenly and resolve better than structured meshes derived from hex elements

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# Taylor Bar: Force and Length History

Convergence of axial reaction force and final bar length:



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## Taylor Bar: Zero Pressure Isosurface

- Zero pressure isosurfaces: regions with no volume change
- We are able to resolve these surfaces very smoothly



Pressure and zero pressure iso-contours at final time

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# **Bending Beam**

- Length/Width: 6/1.4m, density: 1.1e3
- Elastic neo-Hookian material (E=1.7e7, ν=0.45)
- Fixed end, driven by initial x-velocity profile



#### Versions of initial velocity and geometry

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## Bending Beam: Pressure



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# Bending Beam: Force History

- We run on four uniform (unstructured) tet meshes  $(m_0-m_3)$
- Convergence of reaction forces (*x*, *y*) at fixed surface:



# Billet in Compression: Pressure

- Length/Radius: 1.5/1.0cm, density: 1e5
- Elastic-plastic material (E=384.62e9,  $\nu$ =0.423,  $\sigma_y$ =1e9, H=3e9)
- Quasistatic approx. using dynamics (fictitious density, velocity)
- Top: dirichlet uniform velocity, bottom: zero normal displacement



Pressure contours for three meshes (25% compression)

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#### Bar in Tension: Pressure

- Length/Radius: 5.33/0.641cm, density: 1e5
- Elastic-plastic material (*E*=80.2e9, ν=0.29, σ<sub>ν</sub>=0.45e9, *H*=0.13e9)
- Quasistatic approx. using dynamics (fictitious density, velocity)
- Top: zero normal displacement, bottom: dirichlet uniform velocity



Pressure contours for three meshes (0.4cm extension)

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# Impact Test in Complex Geometry

- Length/Radius: 32/3.24mm, density: 8930
- Elastic-plastic material (E=117.0e9,  $\nu$ =0.35,  $\sigma_y$ =0.4e9, H=0.1e9)
- Zero normal velocity at wall, initial uniform x-velocity 100m/s



#### Initial Geometry

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#### Impact Test: Meshes



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#### Impact Test: Meshes



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#### Impact Test: Meshes



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# Impact Test: Deformation & Pressure



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# Impact Test: Deformation & Pressure



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Image: A mathematical states in the second states in the second

# Impact Test: Deformation & Pressure



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Image: A matrix and a matrix

# Impact Test: Velocity & Pressure

Plot on half-domain:



#### Velocity magnitude



#### Pressure

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# Impact Test: Velocity & Pressure

Plot on half-domain:



#### Velocity magnitude



Pressure

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# Impact Test: Velocity & Pressure

Plot on half-domain:



#### Velocity magnitude



Pressure

#### Impact Test: Zero Pressure Isosurfaces



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#### Impact Test: Zero Pressure Isosurfaces



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#### Impact Test: Zero Pressure Isosurfaces



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# Impact Test: Plastic Strain



#### Plastic strain (left) and subgrid scale velocity (right)

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# Summary and Ongoing Work

#### Current status

- Finite deformation solid mechanics capability for tet meshes
- Method is stable and accurate (based on VMS)
- Compatible with VMS-based nodal hydrocode (we have a separate fluid/solid coupling module)

#### Ongoing work

- Additional formal code verification
- Performance improvements
- Comparisons with hex-based solid mechanics codes
- Publications on solids and fluid/solid coupling

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